



Critical advances in charge separation regulation for photocatalytic CO₂ reduction

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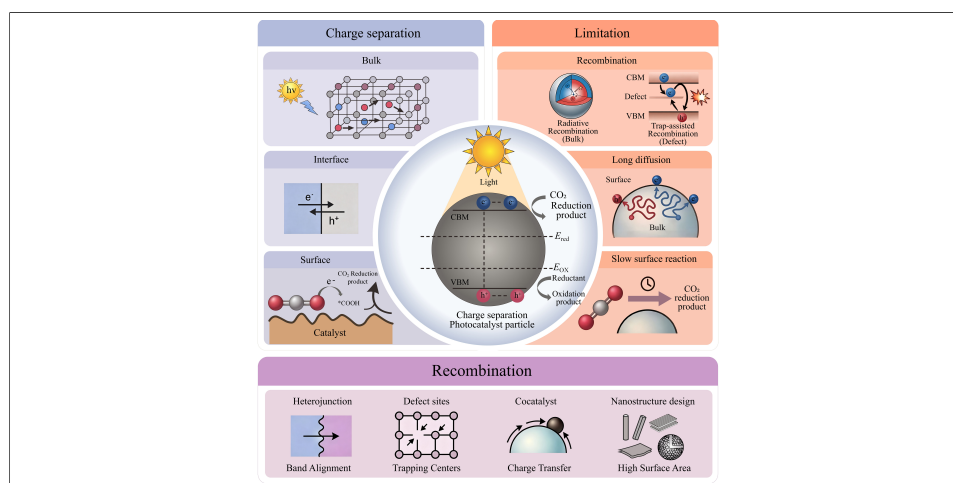
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Abstract

The photocatalytic CO₂ reduction (PCR) technology mimics the energy conversion process of natural photosynthesis. PCR is an important strategy to help achieve the future “dual carbon” goals due to its significant advantages of being green, clean, and sustainable. However, the current PCR technology still has a large gap from practical industrial application owing to key problems such as low CO₂ conversion efficiency, poor product selectivity, and insufficient photocatalyst stability. The fundamental reason is the low charge separation efficiency, which is also the core bottleneck restricting performance improvement. This review first outlines the basic mechanism of PCR and systematically analyzes the physical essence of charge separation processes at spatial scales. It further sorts out the internal and external factors that limit the charge separation efficiency, as well as the common techniques and evaluation standards for characterizing charge behavior. It also highlights representative progress in boosting charge separation efficiency through



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such strategies as heterojunction construction, defect engineering, catalyst regulation, and morphology optimization. This review provides references and inspiration for the rational design of high-performance PCR catalysts.

INTRODUCTION

Since the Industrial Revolution, the excessive use of conventional fossil fuels (coal, oil, and natural gas) has resulted in a continuous increase in CO₂ concentration, which has become a serious challenge for human society^[1,2]. If solar energy can be directly used to convert CO₂ into CO, CH₄, CH₃OH, and even C₂+ high-value-added hydrocarbons and oxygen-containing chemicals^[3], it will be possible to establish a closed loop of “emission reduction-carbon fixation-fuel/chemical production.” This would partially replace fossil resources and alleviate the pressure of carbon emissions. Currently, there are numerous methods for CO₂ reduction, including photocatalysis, electrocatalysis, photoelectrocatalysis, biocatalysis, thermal catalysis, and plasma catalysis^[1]. Among these methods, photocatalytic CO₂ reduction (PCR) mimics natural photosynthesis by using semiconductor photocatalysts to capture solar energy and drive the direct conversion of CO₂ and H₂O into hydrocarbons^[4]. This process relies solely on solar energy as the input and water and CO₂ as feedstocks. Therefore, compared with other CO₂ reduction methods, PCR has prominent advantages of being green, clean, and sustainable. In recent years, with the rapid advancement of valuation standardization and mechanism investigation, the PCR research paradigm has gradually shifted from “material screening” to “mechanism-driven directed design”, and has achieved excellent performance.

However, due to some key challenges and drawbacks, PCR technology is currently at the stage of basic research rather than practical application. For instance, under low CO₂ concentrations, mass transfer and surface adsorption efficiencies are poor^[5]. Furthermore, the inherent chemical inertness of the CO₂ molecule imposes high energy barriers for its activation and subsequent multi-step proton-coupled electron transfer processes^[5]. Lastly, overall photon utilization and quantum efficiencies remain low^[5]. The fundamental cause of these apparent problems is the limited photogenerated charge separation in semiconductor photocatalysts, which remains the core bottleneck in the field of PCR^[6]. This is manifested in the following two aspects. Firstly, photogenerated electron-hole pairs undergo rapid recombination. After the semiconductor photocatalyst absorbs photons, electron-hole pairs are generated, but over 90% of them are typically subjected to non-radiative recombination within picoseconds to nanoseconds through bulk or surface defects^[7], thus failing to participate in the surface catalytic reaction. This causes a significant loss of photogenerated charges. Secondly, the charge transfer resistance is high. For the separated charges that do migrate to the active sites on the catalyst surface, they often face high interface energy barriers, long transfer paths, and slow coupling rates with reactants or intermediates^[8], resulting in sluggish charge utilization kinetics. These bottlenecks not only directly reduce CO₂ conversion activity and stability, but also alter the competitive kinetics of the reaction pathways by influencing the formation and accumulation of key intermediates (such as *CO₂⁻, *COOH, *CO, * and CHO). Therefore, insufficient charge separation of photocatalysts is a key factor restricting the activity, selectivity, and stability of PCR. Addressing this challenge, effective charge separation regulation has become a central focus in research on photocatalytic CO₂ reduction.

To overcome this issue, researchers have proposed various regulation strategies to enhance charge separation and achieve directional transport. There are three typical examples. (1) At the bulk phase level, the charge transfer path is shortened to inhibit the bulk recombination via band engineering, defect/doping regulation, and morphology structure optimization^[9]. (2) At the interface level, spatial charge separation is achieved by constructing various heterojunctions (type-II, Z-scheme, S-scheme, *etc.*)^[2]. (3) At the surface level, the construction of Schottky junctions, loading of cocatalysts, and design of atomic-scale active sites are

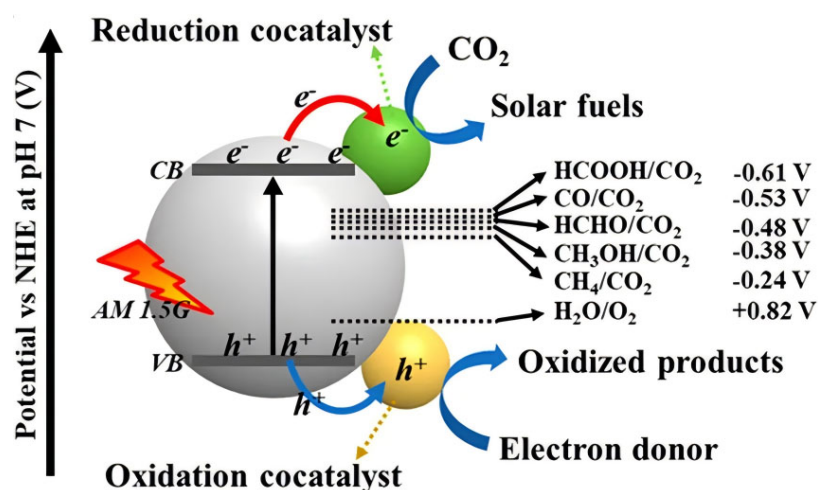


Figure 1. CO₂ conversion by semiconductor photocatalyst. Reprinted from Ref.^[11] under the CC BY license. NHE: Normal hydrogen electrode.

employed to achieve efficient coupling between electron extraction and reaction sites, thereby accelerating the adsorption and activation of CO₂ and lowering the energy barrier of rate-determining reaction steps^[10]. Especially in recent years, the heterojunction and multi-interface coupling systems emphasize enhancing charge separation while retaining strong oxidation/reduction capabilities, providing a possibility for CO₂ multi-electron deep reduction and C-C coupling.

This review provides a systematic overview of charge separation mechanisms in PCR systems. We first discuss the fundamental principles governing photocatalytic CO₂ reduction and charge carrier dynamics, followed by a detailed analysis of intrinsic and extrinsic factors that limit charge separation efficiency. Subsequently, we summarize representative characterization techniques for probing charge behavior across bulk, interface, and surface scales. Finally, we highlight recent advances in charge separation regulation strategies, including heterojunction engineering, defect modulation, cocatalyst optimization, and morphology design, and provide perspectives on future directions for the rational design of high-performance photocatalysts.

FUNDAMENTAL PRINCIPLES OF CHARGE SEPARATION

Principles of photocatalytic CO₂ reduction

The photocatalytic CO₂ reduction process is governed by three sequential steps: light absorption, charge carrier dynamics, and surface redox reactions, as illustrated in Figure 1. Upon irradiation, the semiconductor photocatalyst absorbs photons with energy equal to or greater than its band gap ($h\nu \geq E_g$). This results in the excitation of electrons from the valence band (VB) to the conduction band (CB), thereby generating electron-hole pairs.

The photogenerated electrons in the CB migrate to the catalyst surface and can be further transferred to a reduction cocatalyst. Here, they participate in the multi-electron reduction of adsorbed CO₂ into solar fuels such as HCOOH, CO, HCHO, CH₃OH, and CH₄^[4]. Thermodynamically, this requires the CB minimum to be more negative than the corresponding CO₂ reduction potentials. Meanwhile, the photogenerated holes remaining in the VB migrate to an oxidation cocatalyst to oxidize electron donors (e.g., H₂O or sacrificial agents)^[6] into oxidized products, provided that the VB potential is more positive than the oxidation potential.

Therefore, the overall efficiency of photocatalytic CO₂ reduction is synergistically determined by the light absorption capacity, the separation and migration efficiency of photogenerated charge carriers, and the kinetics of surface redox reactions. In particular, efficient charge separation and transfer are critical, as rapid electron-hole recombination during migration severely limits catalytic performance. Consequently, suppressing charge recombination and promoting directional charge transfer are paramount strategies for enhancing the overall PCR efficiency.

Charge separation process

Bulk charge separation

Spatially, charge separation encompasses bulk diffusion to the semiconductor surface and interfacial transfer between different components in composites^[12]. In a pristine photocatalyst, the bulk carrier diffusion length critically dictates the efficiency of charge separation and subsequent participation in CO₂ reduction^[13]. If the bulk diffusion path is too long, photogenerated electrons will suffer from severe bulk recombination before reaching the active surface sites, rendering even strong light absorption ineffective. Conversely, minimizing the diffusion distance facilitates rapid carrier migration to the active sites. Consequently, engineering low-dimensional nanostructures or hierarchical porous architectures is a highly effective strategy for optimizing bulk charge separation.

However, the performance of bare photocatalysts remains constrained by rapid bulk charge recombination. To mitigate this, composite photocatalysts are constructed by integrating two semiconductors with appropriately matched energy bands^[7]. Based on their relative band alignments, heterojunctions are generally classified into three configurations [Figure 2A]: type-I (straddling gap), type-II (staggered gap), and type-III (broken gap)^[14]. In a type-I alignment, both electrons and holes transfer to the narrower-bandgap semiconductor, failing to achieve spatial charge separation. Meanwhile, for a type-III alignment, the extreme band offset renders continuous charge transfer energetically prohibitive. In a type-II heterojunction, electrons migrate to the lower conduction band while holes transfer to the higher valence band. This staggered transfer effectively separates the electron-hole pairs spatially, thereby enhancing photocatalytic performance^[2]. However, because electrons accumulate at a less negative conduction band potential, the system's overall reductive driving force is compromised, which significantly limits multi-electron CO₂ reduction kinetics. Alternatively, staggered band structures can facilitate a completely different charge transfer mechanism known as a Z-scheme system, where the charge flow resembles the letter "Z" [Figure 2B]. Depending on the charge-mediating mechanism, Z-scheme systems are classified into traditional (using a redox mediator), all-solid-state (using a solid conductor), and direct Z-scheme (without an intermediate medium) heterojunctions. Compared with traditional Z-scheme and all-solid-state Z-scheme, the direct Z-scheme heterojunction can avoid the selection and shielding effects of electron media and reverse photochemical reactions^[14]. In this unique pathway, the photogenerated electrons with weaker reductive capability preferentially recombine with holes possessing weaker oxidative capability. Consequently, the strongly reductive electrons and strongly oxidative holes are preserved in their respective conduction and valence bands. Thus, the direct Z-scheme configuration successfully circumvents the redox potential loss inherent to type-II systems. Therefore, constructing appropriate heterojunctions is a critical strategy for enhancing interfacial and spatial charge separation. XPS is an important means for determining the chemical composition and electronic states of materials^[15].

Interface charge separation

Photogenerated charge separation at the interface of composite materials is a critical determinant of both the efficiency and selectivity of CO₂ reduction. A prominent example is the charge transfer mechanism in S-scheme heterojunctions, a concept initially introduced by Xu *et al.* in 2019^[15]. An S-scheme system consists of a reducing photocatalyst (RP) and an oxidizing photocatalyst (OP). Upon contact, the difference in their

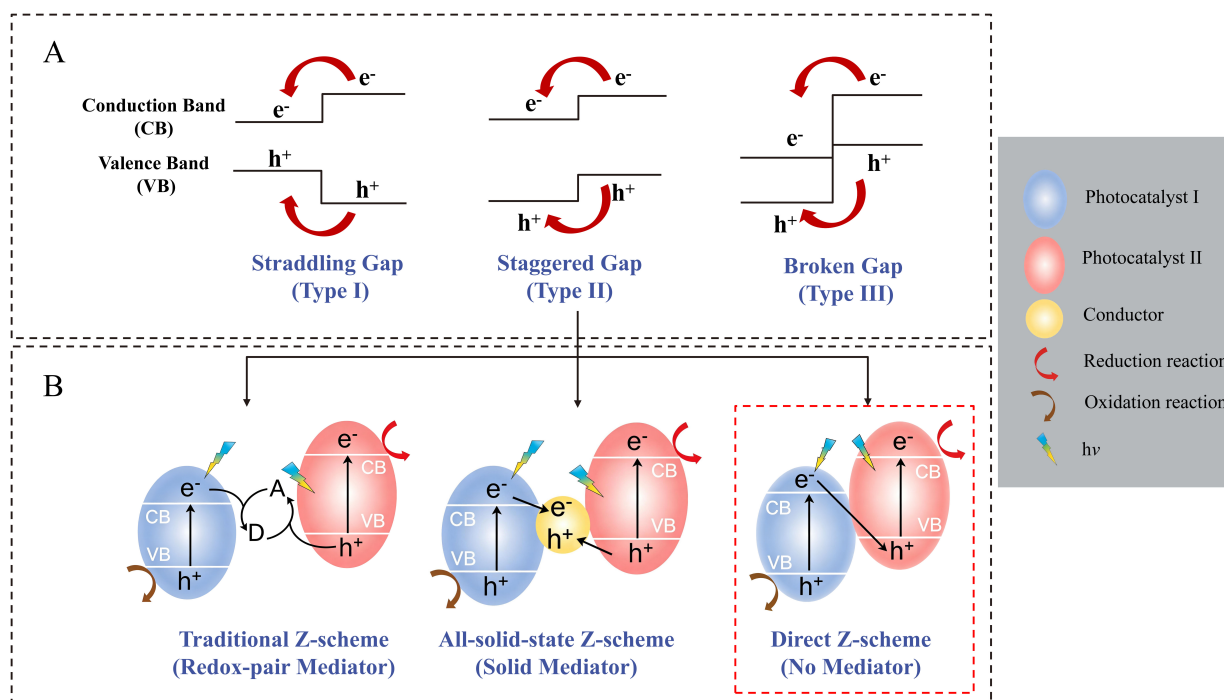


Figure 2. Schematic illustration of typical charge separation paths in two-semiconductor photocatalysts. (A) Band alignment; (B) Z-scheme heterojunctions. Figure 2A and B are adapted from Ref.^[14] under the CC BY-NC 3.0 license.

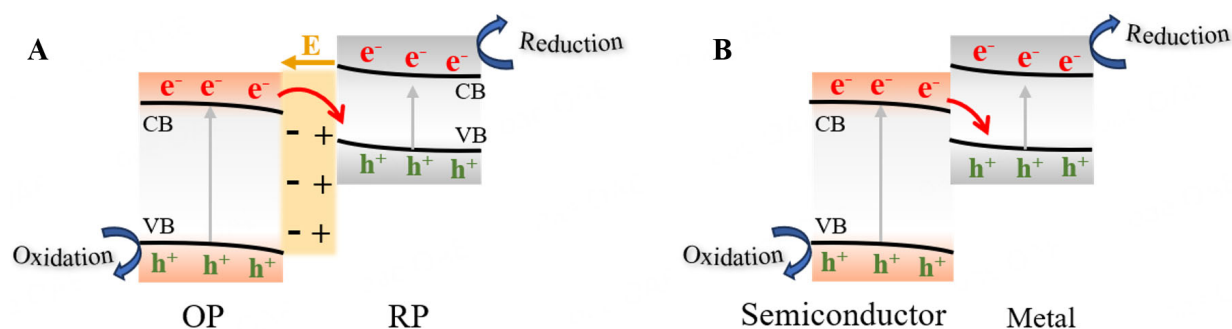


Figure 3. Schematic illustration of typical charge separation paths: (A) S-scheme heterojunction. E is the built-in electric field. (B) Schottky heterojunction. Created with Microsoft PowerPoint. CB: Conduction band; VB: valence band; OP: oxidizing photocatalyst; RP: reducing photocatalyst.

Fermi levels drives spontaneous electron transfer, establishing an internal electric field and inducing corresponding band bending at the interface. Under illumination, this built-in field, synergized with band bending, drives a macroscopic “step-like” interfacial charge transfer pathway [Figure 3A]. This mechanism selectively facilitates the recombination of the less reductive electrons in the OP and the less oxidative holes in the RP. Consequently, the strongly reductive electrons and strongly oxidative holes are preserved in their respective conduction and valence bands. Thus, interfacial charge separation in S-scheme systems is highly directional, achieving efficient carrier separation while maximizing the overall redox potential.

Furthermore, even if photogenerated carriers successfully migrate from the bulk to the semiconductor surface, they remain highly susceptible to rapid recombination if not properly extracted. Therefore, electrons must be further directed across the solid-liquid interface. To achieve this, metallic materials are typically deposited onto the semiconductor surface to form a Schottky junction [Figure 3B]^[16]. Because the metal generally possesses a larger work function, electrons at the semiconductor surface are preferentially injected

into the metal, which acts as an effective electron trap. Ultimately, these electrons are successfully funneled to the solid-liquid interface, where they participate in the CO₂ reduction reaction on the catalyst surface. The formation of a Schottky barrier subsequently prevents electron backflow, ensuring the strict directionality of surface charge transfer.

Thermodynamic and kinetic parameters of charge separation

In photocatalytic CO₂ reduction, charge separation efficiency is jointly governed by thermodynamic driving forces and kinetic competition^[17]. Thermodynamically, this depends on the band alignment between the semiconductor and the CO₂ reduction potentials, which must provide sufficient driving force for the redox reactions while minimizing excessive energy loss^[1]. The kinetic competition is primarily determined by the relative charge transfer rate (including interfacial charge injection and transport) and the charge recombination rate (non-radiative recombination in the bulk, surface, and interface)^[18]. When the charge transfer rate is significantly higher than the recombination rate, the charge separation efficiency is effectively enhanced. The charge separation efficiency can be expressed as^[18]: $\eta_{\text{sep}} = \frac{K_{\text{ct}}}{K_{\text{ct}} + K_{\text{rec}}}$, where η_{sep} represents the charge separation efficiency, defined as the fraction of photogenerated electron-hole pairs that are successfully separated and subsequently participate in catalytic reactions. A higher η indicates a greater population of charge carriers effectively involved in catalysis, thereby corresponding to enhanced photocatalytic performance. This parameter fundamentally reflects the dynamic competition between charge transfer and recombination processes. The parameter K_{ct} denotes the charge transfer rate constant, which describes the kinetics of photogenerated carrier migration from the semiconductor bulk or interface to active reaction sites. This encompasses processes such as bulk-to-surface diffusion, carrier injection across semiconductor/cocatalyst interfaces, and electron transfer to adsorbed CO₂ molecules. A larger K_{ct} implies more efficient carrier delivery to catalytic sites, thereby facilitating reaction kinetics. In contrast, K_{rec} represents the charge recombination rate constant, describing the rate at which electrons and holes recombine, including bulk recombination, surface recombination, and interfacial recombination. An increased K_{rec} indicates more severe carrier loss, which is detrimental to photocatalytic efficiency. When K_{ct} is much higher than K_{rec} , the photogenerated charges maintain an effective lifetime in a scale of nanosecond to microsecond. This prolonged lifetime provides a sustained electron supply for the continuous generation of key reactive intermediates (e.g., *CO₂ and *COOH). Furthermore, localized internal electric fields, Fermi level alignments, and adsorption-induced dipoles can accelerate interfacial charge injection. These factors enable the rapid transfer of electrons to metal, single-atom, or carbon-based active sites, while holes are concurrently scavenged by strong oxidation centers^[17]. This synergistic effect intrinsically suppresses recombination and drastically enhances the overall charge separation efficiency.

Key factors limiting charge separation

The efficiency of charge separation in photocatalytic CO₂ reduction (PCR) is dictated by a combination of both intrinsic and extrinsic factors, which collectively regulate carrier generation, migration, and utilization across the bulk, interface, and surface of photocatalysts^[7]. Distinguishing between these two categories is crucial for understanding their respective roles and interdependent effects^[17].

Intrinsic factors are primarily dictated by the physicochemical properties of the photocatalyst. At the bulk level, rapid recombination of photogenerated electron-hole pairs represents the primary source of energy loss. In narrow-bandgap semiconductors, high intrinsic carrier densities can exacerbate Coulombic attraction, leading to severe bulk recombination. In addition, deep-level defects and structural disorder introduce trap states that act as non-radiative recombination centers, accelerating carrier annihilation during migration^[16]. At the interface level, intrinsic limitations arise from unfavorable band alignment and lattice mismatch. Improper energy level matching between semiconductors or between semiconductors and cocatalysts impedes the formation of a robust built-in electric field, thereby weakening the driving force for

directional charge transfer. Consequently, carriers often recombine at the interface before completing migration^[18]. At the surface level, insufficient exposure of active sites or mismatched electronic structures (e.g., inadequate electron affinity) can impede charge utilization. Electrons failing to rapidly inject into adsorbed CO₂ molecules remain highly susceptible to recombination. Concurrently, holes that are not efficiently scavenged by oxidation reactions accumulate at surface states, inducing additional recombination losses^[18]. Collectively, these intrinsic constraints severely limit the fraction of charge carriers available for catalytic reactions^[10].

In contrast, extrinsic factors stem from external operating conditions and the reaction environment. A primary constraint is restricted light-harvesting capability, particularly in wide-bandgap oxide photocatalysts that respond primarily to ultraviolet light^[1]. This limited photon absorption not only curtails the initial generation of charge carriers but also creates a spatial mismatch between the regions of carrier generation and reaction sites, increasing the probability of recombination during migration.

Furthermore, sluggish activation and weak adsorption of CO₂ molecules significantly impede charge utilization. When CO₂ concentration is low or the catalyst-CO₂ interaction is weak, sluggish activation and weak adsorption of CO₂ molecules significantly impede charge utilization (e.g., *CO₂⁻)^[5]. Consequently, accumulated electrons at the interface are more likely to recombine with holes rather than participate in reduction reactions. Because CO₂ reduction is a complex multi-electron transfer process, any disruption can ultimately degrade catalytic activity, product selectivity, and apparent quantum efficiency^[1].

Ultimately, intrinsic and extrinsic factors do not operate in isolation; rather, they compound one another to dictate the overall charge separation efficiency. Therefore, effective regulation strategies should simultaneously address material design and reaction conditions to minimize recombination losses and maximize carrier utilization^[7].

Characterization techniques for charge behavior

To elucidate the intrinsic mechanisms of charge separation in photocatalytic CO₂ reduction (PCR) systems, it is essential to transcend isolated descriptors and instead consider charge behavior as a continuous physicochemical process spanning photogeneration, carrier migration^[19], interfacial redistribution, and surface reaction coupling, as schematically integrated in Figure 4. In this context, characterization techniques should be interpreted in direct correspondence with these successive stages, rather than as independent analytical tools, so as to establish a mechanistically consistent understanding of charge evolution.

At the initial stage of light absorption and carrier generation (Step 1-2 in Figure 4), the dominant process is the competition between carrier excitation and ultrafast recombination in the bulk phase. The absorption characteristics of light can be analyzed through ultraviolet–visible (UV-vis) diffuse reflectance spectroscopy, and the band gap (E_g) value of the photocatalyst can be estimated using the Tauc plot method^[20]. Photoluminescence (PL) and time-resolved photoluminescence (TRPL) are therefore commonly employed to probe recombination dynamics. An attenuated PL intensity or prolonged lifetime is often interpreted as indicative of reduced recombination probability^[21]. However, such interpretations are inherently limited because PL-based techniques selectively probe radiative recombination pathways and cannot distinguish whether the attenuation of emission arises from improved spatial charge separation or enhanced non-radiative decay via defect states^[22]. In addition, TRPL provides only averaged lifetime information and lacks the spatial resolution required to differentiate recombination occurring in the bulk, at the interface, or on the surface. Consequently, although PL and TRPL are useful for preliminary screening, they cannot independently establish the effectiveness of charge separation along the migration pathway depicted in Figure 4. For instance, in a Co₁-C₃N₄@ α -Fe₂O₃ Z-scheme photocatalyst, TRPL analysis revealed a prolonged

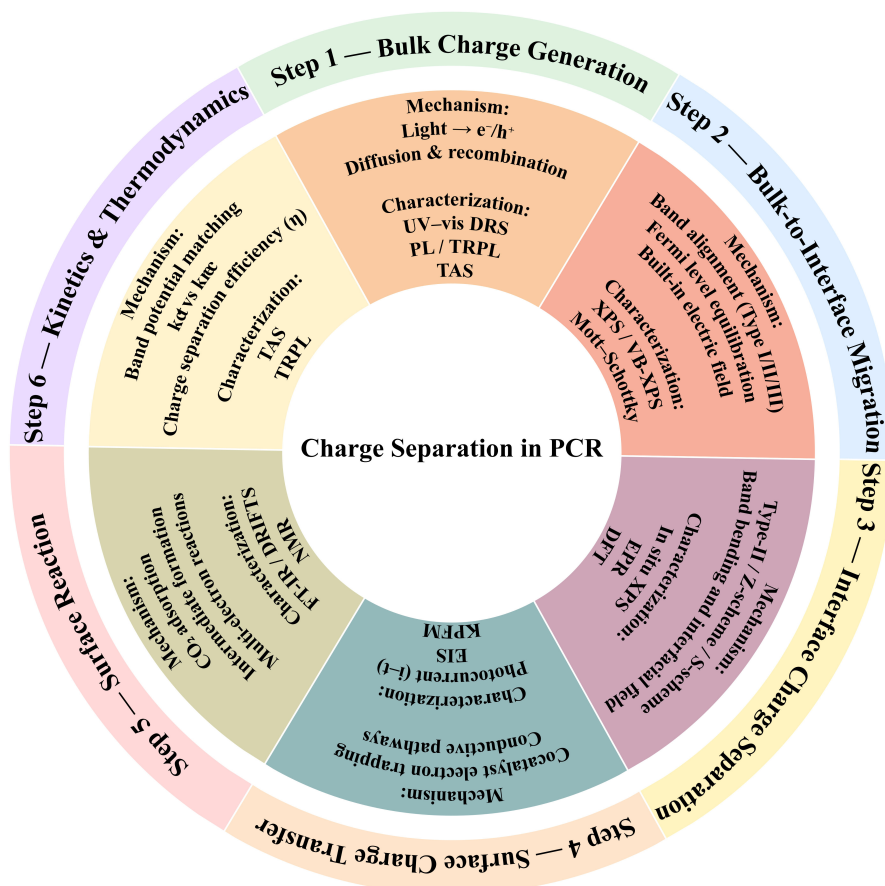


Figure 4. Integrated correlation between charge separation processes and characterization techniques across bulk, interface, and surface levels in PCR systems. Created with Microsoft PowerPoint. PCR: Photocatalytic CO_2 reduction; TAS: transient absorption spectroscopy; TRPL: time-resolved photoluminescence; NMR: nuclear magnetic resonance; FT-IR: Fourier transform infrared spectroscopy; EIS: electrochemical impedance spectroscopy; KPFM: Kelvin probe force microscopy; EPR: electron paramagnetic resonance; DFT: density functional theory; XPS: X-ray photoelectron spectroscopy; VB: valence band; UV-vis: ultraviolet–visible; DRS: diffuse reflectance spectroscopy; PL: photoluminescence.

exciton lifetime after constructing the heterojunction and introducing single-atomic Co sites, indicating that interfacial charge separation and atomic-site regulation synergistically attenuated carrier recombination^[23]. This example demonstrates that PL/TRPL can provide preliminary evidence for recombination inhibition, although such signals should still be correlated with non-radiative pathways and spatial charge migration.

As photogenerated carriers migrate from the bulk toward interfaces (Step 2-3 in Figure 4), the efficiency of charge extraction and transport becomes the critical factor governing subsequent reaction steps. Electrochemical techniques, such as transient photocurrent response and electrochemical impedance spectroscopy (EIS), are frequently used to evaluate this process^[24]. Enhanced photocurrent is generally associated with improved carrier mobility and extraction efficiency, while reduced charge-transfer resistance inferred from Nyquist plots is often taken as evidence of facilitated interfacial transport. Nevertheless, these electrochemical responses are intrinsically influenced by extrinsic factors, including electrode configuration, electrolyte diffusion, and interfacial capacitance, which complicates their direct correlation with intrinsic charge-transfer pathways. As a result, EIS cannot unambiguously resolve whether the observed improvement originates from genuine directional charge separation or from changes in macroscopic transport properties, and therefore cannot distinguish between different charge-transfer modes embedded in Figure 4, such as type-II versus S-scheme mechanisms^[25]. In a $Cu_2O@Cu$ -TCPP heterojunction for CO_2 photoreduction, the

optimized Cu₂O@CP-1 sample exhibited the strongest transient photocurrent response and the smallest EIS semicircle, confirming that the *in situ*-constructed p-n junction effectively promoted charge separation and interfacial transport. This case illustrates that TPC and EIS are useful macroscopic indicators of carrier extraction efficiency and charge-transfer resistance.

To directly capture the dynamic evolution of carriers along the migration-separation pathway (Step 2-4 in Figure 4), transient absorption spectroscopy (TAS) provides a uniquely powerful approach by resolving carrier kinetics across femtosecond-to-microsecond timescales^[21]. By tracking the temporal evolution of excited-state species, TAS enables direct observation of carrier trapping, relaxation, and interfacial transfer processes that are otherwise inaccessible by steady-state techniques. For instance, Wang *et al.* employed femtosecond transient absorption spectroscopy, *in situ* irradiated X-ray photoelectron spectroscopy, and electron paramagnetic resonance to investigate a fibrous Ta₂O₅/Ag₂S S-scheme photocatalyst for diluted CO₂ photoreduction^[21]. The TAS results revealed ultrafast decay components corresponding to the selective recombination of low-energy carriers, while long-lived signals were assigned to high-energy electrons retained for reduction reactions, providing direct kinetic evidence for the selective recombination-retention mechanism characteristic of S-scheme charge transfer^[26].

While TAS resolves the temporal dimension of charge evolution, direct identification of charge redistribution at heterointerfaces (Step 3-4 in Figure 4) requires operando-level electronic structure characterization. In this regard, *in situ* irradiated X-ray photoelectron spectroscopy (ISIXPS) has emerged as one of the most compelling techniques for validating charge-transfer directionality^[27]. Under illumination, shifts in core-level binding energies directly reflect electron accumulation or depletion in different components of a heterojunction, thereby providing real-time evidence of charge migration across interfaces^[28,29]. A landmark example is the TiO₂@ZnIn₂S₄ photocatalyst reported by Wang *et al.*^[28], in which opposite photoinduced binding-energy shifts were observed for the two semiconductors under irradiation, unambiguously confirming directional electron transfer driven by the built-in electric field. This *operando* evidence directly maps onto the interfacial charge-separation process depicted in Figure 4 and has become one of the most widely accepted experimental criteria for distinguishing S-scheme mechanisms from conventional type-II charge-transfer models^[28].

At the final stage, where separated charges participate in surface reactions (Step 4-6 in Figure 4), the effectiveness of charge separation must ultimately be evaluated in terms of its coupling with CO₂ adsorption and activation. Kelvin probe force microscopy (KPFM) provides nanoscale mapping of surface potential and work-function distribution, enabling visualization of spatial charge accumulation at active sites and interfaces under illumination^[30]. At the nanoscale, KPFM has been used to directly visualize photo-injected electrons in CdS quantum dot-modified cesium tungstate nanosheets. The light-induced surface-potential variation revealed directional electron injection from the photoexcited CdS quantum dots into the metal oxide nanosheets, thereby mapping spatial charge redistribution at the heterointerface^[31]. This example demonstrates that KPFM can correlate local work-function changes with charge accumulation regions, providing spatial evidence complementary to spectroscopic and electrochemical analyses. Meanwhile, *in situ* spectroscopic techniques such as DRIFTS and electron paramagnetic resonance (EPR) directly track the formation and evolution of reaction intermediates, including *CO₂⁻ and *COOH species, thereby linking charge injection to specific elementary reaction steps^[32]. These approaches highlight that efficient charge separation, as outlined in Figure 4, only translates into enhanced photocatalytic performance when it is effectively coupled with surface reaction kinetics. At the reaction level, *in situ* DRIFTS and EPR can further connect separated charges with CO₂ activation. In CsBr@CuBr₂ and related halogen-defect systems, *in situ* DRIFTS was used to follow the evolution of CO₂-derived intermediates, while EPR-related analysis revealed light-induced defect states that acted as dynamic active sites for CO₂ photoreduction^[33]. These observations

establish a mechanistic link between charge trapping, defect-mediated CO₂ activation, and intermediate evolution, thereby providing reaction-level evidence for charge utilization.

Despite the apparent diversity of characterization methods, it is crucial to recognize that each technique probes only a specific segment of the charge evolution pathway illustrated in Figure 4. PL and TRPL primarily reflect recombination behavior in the early stage, electrochemical methods capture macroscopic transport properties during migration, TAS resolves ultrafast carrier kinetics, whereas ISIXPS and KPFM provide direct information on interfacial and spatial charge redistribution. Therefore, reliable identification of charge-transfer mechanisms in PCR systems requires convergent evidence from multiple complementary techniques that collectively reconstruct the full charge separation process from generation to reaction. Recent methodological analyses have further emphasized that operando ISIXPS, ultrafast spectroscopy, and theoretical calculations should be jointly employed to establish a self-consistent charge-transfer framework^[27].

Accordingly, the characterization techniques discussed above should not merely be regarded as diagnostic tools but as essential criteria for evaluating the validity of proposed charge-separation models. In the following sections, these methods will be revisited in representative case studies to correlate experimentally validated charge-transfer pathways with photocatalytic performance, thereby establishing a unified framework linking Figure 4 to rational catalyst design.

CHARGE SEPARATION REGULATION STRATEGIES

Heterojunction construction

Heterojunction engineering stands as a paramount strategy for regulating charge separation and directional carrier transport in photocatalytic CO₂ reduction systems. From the perspectives of charge-transfer pathways and band alignment, heterojunctions can be generally categorized into conventional type-II, advanced systems (Z-scheme and S-scheme), and Schottky junctions^[34].

Historically, the type-II heterojunction has served as the classical model for achieving spatial separation of photogenerated electrons and holes through staggered band alignment. While this configuration effectively suppresses carrier recombination, it inevitably drives electrons and holes toward lower-energy band edges, precipitating a substantial loss of reduction and oxidation potentials^[35]. This inherent trade-off fundamentally limits its applicability in photocatalytic CO₂ reduction, where multi-electron transfer reactions demand robust thermodynamic driving forces.

Consequently, recent studies have progressively shifted from conventional type-II systems toward advanced heterojunction configurations that can simultaneously promote charge separation while preserving high-energy carriers. In this context, Z-scheme and S-scheme heterojunctions achieve the selective recombination of low-energy carriers and the retention of strong redox potentials through distinct interfacial charge-transfer pathways. This represents a paradigm evolution from “separation-dominated” to “energy-preserved” charge regulation. In parallel, Schottky heterojunctions, formed at semiconductor-metal interfaces, provide an alternative route by enabling efficient electron extraction and directional transfer via built-in Schottky barriers. This curtails surface recombination without relying on inter-semiconductor band alignment.

Guided by these considerations, the following sections will focus on Z-scheme, S-scheme, and Schottky heterojunctions as representative advanced models for efficient charge separation and utilization. Meanwhile, the conventional type-II mechanism will serve primarily as a foundational reference framework for understanding the evolution of heterojunction design.

Z-scheme heterojunction

The Z-scheme heterojunction is derived from the “Z-scheme” electron transfer mechanism of natural photosynthesis^[36]. Its core advantage lies in achieving efficient separation of photogenerated electrons and holes while retaining high-energy electrons with strong reducing ability and high-energy holes with strong oxidizing ability to the greatest extent^[37]. This not only provides sufficient high-energy electron donors for the CO₂ photoreduction reaction but also lays a thermodynamic foundation for multi-step proton-coupled electron transfer (PCET) processes, making it an important design direction for catalytic systems to enhance the performance of CO₂ photoreduction.

Fundamentally, the direct Z-scheme heterojunction is an effective charge-transfer regulation strategy^[38], and its fundamental charge transfer mechanism is illustrated in Figure 5A. A typical Z-scheme system couples a semiconductor possessing strong oxidation capabilities (analogous to Photosystem II) with another possessing strong reduction capabilities (analogous to Photosystem I). After the two come into contact, the low-energy electrons in the conduction band of PS I and the low-energy holes in the valence band of PS II will preferentially recombine and be consumed at the interface; ultimately the high-energy electrons with strong reduction ability on the conduction band of PS II (directly used for the activation and reduction of CO₂), and the high-energy holes with strong oxidation ability on the valence band of PS I (used to oxidize sacrificial agents) are retained. This characteristic recombination-retention pathway enables both efficient charge separation and strong redox capability. As shown in Figure 5A, the photogenerated electrons in semiconductor B are excited to its conduction band, while holes remain in the valence band. Meanwhile, semiconductor A undergoes a similar excitation process. The key feature is that the electrons in the conduction band of semiconductor A recombine with the holes in the valence band of semiconductor B at the interface (indicated by the downward arrow in Figure 5A), forming a typical “recombination-retention” pathway. Consequently, the electrons with strong reduction ability are preserved in the conduction band of semiconductor B, while the holes with strong oxidation ability remain in the valence band of semiconductor A. Compared to conventional type-II heterojunctions, this architecture offers two distinct advantages. First, it enhances directional charge transfer, wherein built-in electric fields, Fermi level equilibration, and interfacial barrier effects synergistically suppress bulk carrier recombination. Second, it preserves thermodynamic driving forces by circumventing the transfer of carriers to weakly redox-active band edges. The retained high-energy electrons have a sufficiently negative chemical potential, which can efficiently overcome the highest energy barrier of the first electron injection step in the CO₂ reduction reaction, and at the same time provide sufficient energy support for subsequent multi-electron coupling processes, ensuring the continuous progress of the reduction reaction. For further reduction of CO₂, especially the further hydrogenation of *CO intermediates or C-C coupling and other complex steps, the value of the Z-scheme heterojunction is not limited to the basic charge separation, but is reflected in three key dimensions: Firstly, the precise retention of high-energy active electrons can provide the power for the further reduction steps with high energy barriers; Secondly, promoting the rapid cross-interface charge transfer can reduce energy loss during the migration process; Thirdly, maintaining the high availability of electrons at the interface ensures that there is always an adequate electron supply at the reduction sites.

In 2020, Wang *et al.*^[41] reported that the Cu₂O-Pt/SiC/IrO_x composite system achieved the coupled photocatalytic reduction of CO₂ to formic acid (HCOOH) and oxidation of water to oxygen (O₂) through the direct and indirect Z-scheme connection of the reduction and oxidation half-reactions. As illustrated in Figure 5D, the Cu₂O-Pt/SiC/IrO_x system integrates spatially separated reduction and oxidation units, in which Cu₂O-Pt mainly serves as the CO₂ reduction side, while SiC/IrO_x functions as the water oxidation side. The Z-scheme charge-transfer pathway enables the recombination of low-energy carriers while retaining electrons with strong reduction ability for CO₂-to-HCOOH conversion and holes with strong oxidation ability for water oxidation. The enhanced charge separation and utilization efficiency are further supported

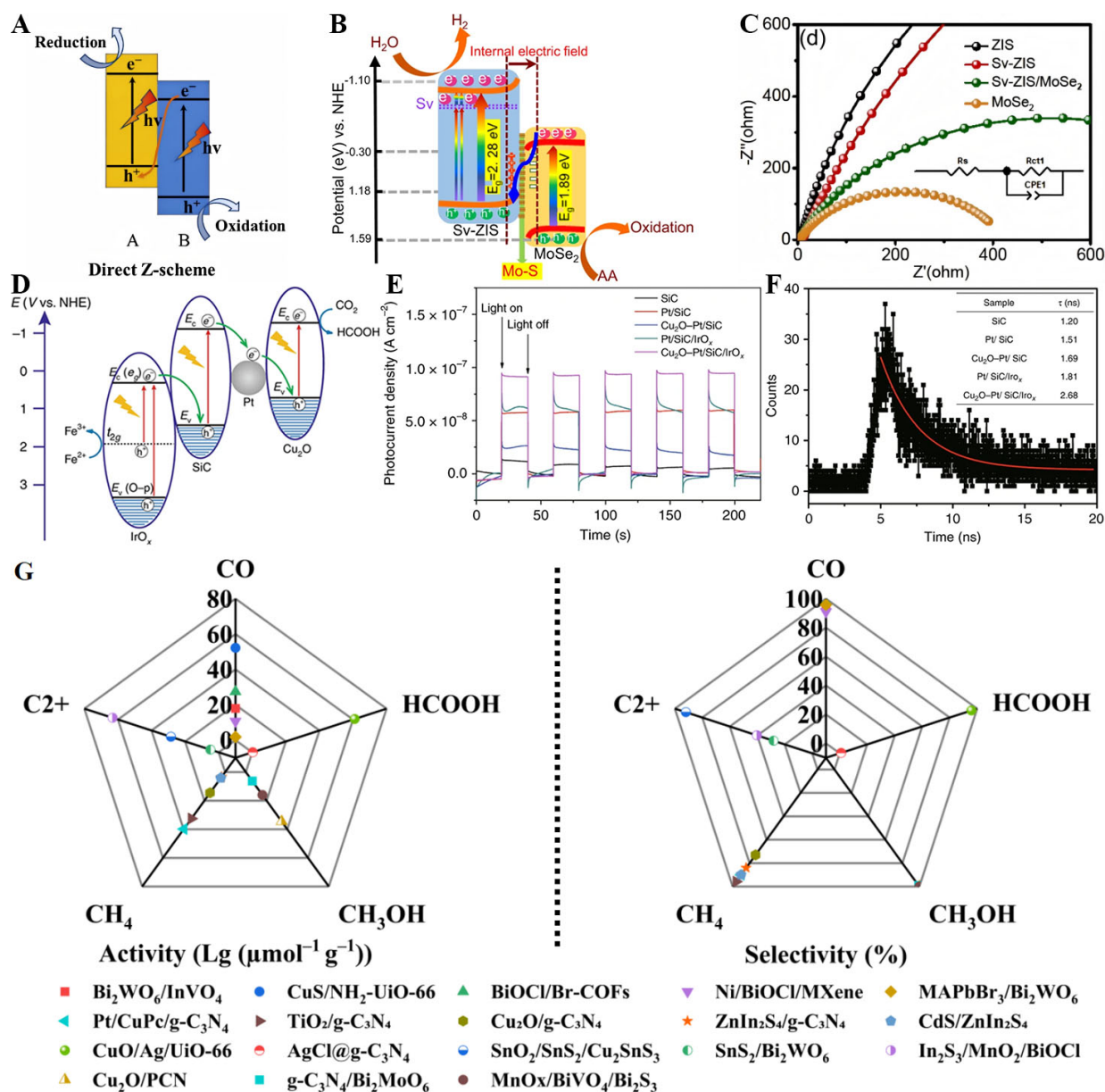


Figure 5. (A) Schematic illustration of charge carrier separation and transfer in a direct Z-scheme heterojunction. Reprinted from Ref.^[39] under the CC BY 4.0 license; (B) Schematic illustration of the Z-scheme charge transfer mechanism mediated by interfacial Mo-S chemical bonds; (C) Charge density difference demonstrating electron redistribution at the interface, indicating the role of Mo-S bonds as fast charge transfer channels. Figure 5B and C is reprinted from Ref.^[40] under the CC BY 4.0 license; (D-F) Charge transfer pathways and corresponding photocatalytic performance in representative Z-scheme heterojunction systems. Figure 5D-F is reprinted from Ref.^[41] under the CC BY 4.0 license; (G) Radar plot comparing catalytic activity (expressed as the logarithm of production rate per mass of catalyst) and selectivity among carbon-containing products. NHE: Normal hydrogen electrode; ZIS: ZnIn₂S₄; AA: ascorbic acid.

by Figure 5E and F. The higher HCOOH and O₂ evolution rates indicate that the photogenerated electrons and holes are effectively consumed in the reduction and oxidation half-reactions, respectively. This balanced production behavior suggests that the Z-scheme configuration promotes efficient charge separation and suppresses the reverse reaction, thereby improving the overall photocatalytic CO₂ reduction performance.

Beyond basic composite designs, the integration of multi-component architectures and interface engineering has further amplified Z-scheme performance. In the ternary ZnFe₂O₄/ZnO/CdS system, the Z-scheme charge transport path was constructed by controlling the interface structure to achieve efficient carrier separation^[42].

The yield of CO_2 to CH_4 was significantly higher than that of single or simple composite systems. The Z-scheme heterojunction constructed through synergy of interface and defect engineering also exhibits remarkable advantages. In the $\text{ZnO}/\text{ZnAl}_2\text{O}_4$ composite, the introduction of oxygen vacancies and interface defects not only improves the interface charge transfer impedance but also promotes the formation of key intermediates (such as $^*\text{COOH}$, $^*\text{HCO}_3^-$), making the conversion performance of CO_2 to CO approximately three times that of pure ZnAl_2O_4 ^[43]. At the same time, both the quantum efficiency and the interface charge separation efficiency have been enhanced. This indicates that the Z-scheme interface has a significant effect in enhancing charge directional migration and reaction kinetics. Recently, some novel Z-scheme heterojunctions have been reported. For instance, in Z-scheme $\text{CuS}/\text{PCN-222}$ heterojunction, an intimate contact and efficient electron-hole separation were achieved at the interface between CuS and PCN-222 , thereby further enhancing the CO_2 reduction activity and carrier utilization rate. This demonstrated the development potential of Z-scheme MOF/sulfide heterojunctions in PCR^[44]. Such studies generally indicate that the charge dynamics in Z-scheme heterojunctions can be significantly optimized by regulating the charge separation pathways through interface structure design (such as interface chemical coupling, defect introduction, and multiphase energy level platforms). Beyond constructing intimate interfacial contact, the introduction of interfacial chemical bonding has emerged as a more effective strategy for regulating charge transfer. The formation of covalent or strongly coupled chemical bonds at heterojunction interfaces can not only significantly reduce interfacial charge transfer resistance, but also provide stable and continuous electron transport pathways, thereby enabling more efficient charge separation and directional migration. Recent studies further reveal that interfacial chemical bonding plays a crucial role in charge transfer regulation. As illustrated in Figure 5B and C, the introduction of interfacial bonds (e.g., Mo-S bonds) can establish strong electronic coupling, optimize band alignment, and facilitate directional charge migration. For instance, in a $\text{Bi}_{19}\text{S}_{27}\text{Br}_3/\text{g-C}_3\text{N}_4$ direct Z-scheme system, the construction of interfacial C-S bonds establishes strong electronic coupling between the two components^[45]. This not only accelerates the transfer of photogenerated electrons from $\text{g-C}_3\text{N}_4$ to $\text{Bi}_{19}\text{S}_{27}\text{Br}_3$, but also increases interfacial electron density and enhances CO_2 adsorption capability. Although the CO_2 adsorption energy of the $\text{Bi}_{19}\text{S}_{27}\text{Br}_3/\text{g-C}_3\text{N}_4$ composite is slightly lower than that of pure $\text{Bi}_{19}\text{S}_{27}\text{Br}_3$, this limitation is effectively compensated by the significantly improved charge separation efficiency arising from the Z-scheme heterojunction and interfacial C-S bonds, ultimately leading to markedly enhanced photocatalytic CO_2 reduction performance^[45]. Charge density difference analysis further confirms that such interfacial bonds act as efficient charge transfer channels [Figure 5F]. Compared to traditional type-II heterojunctions in which electrons/holes separately migrate to lower energy levels and lose their reduction/oxidation capabilities, the Z-scheme heterojunction effectively retains high-energy carriers and reduces bulk-phase recombination through interface advantages^[38], thereby achieving high charge separation efficiency in CO_2 photoreduction reactions and maintaining strong redox driving forces, which is conducive to the complex multi-electron CO_2 reduction (such as the formation of CH_4 , C_2^+ products)^[39].

In addition, Figure 5G summarizes the photocatalytic CO_2 reduction performance of representative Z-scheme heterojunctions by correlating activity and product selectivity for the same material systems. The performance parameters were collected from representative studies reported in Refs.^[46-62] and further reorganized and visualized by the authors to enable a comparative analysis across different Z-scheme heterojunction systems. The left panel compares the production rates of different reduction products, including CO, CH_4 , CH_3OH , and HCOOH , while the right panel presents the corresponding selectivity toward carbon-containing products for each catalyst. This paired comparison shows that Z-scheme heterojunctions not only enhance the overall CO_2 conversion activity but also regulate the product distribution. This unequivocally demonstrates that efficient interfacial charge separation, coupled with the thermodynamic retention of high-energy electrons is the critical determinant in steering the complex reaction pathways of multi-electron CO_2 reduction.

S-scheme heterojunction

In recent years, step-scheme (S-scheme) heterojunctions have rapidly emerged as a focal point in photocatalysis, driven by their unique capacity to regulate interfacial charge transport through a characteristic S-shaped pathway^[15]. Typically composed of two semiconductors with distinct work functions^[54], the S-scheme architecture leverages this work function differential to drive the formation of an intrinsic electric field. As illustrated in [Figure 6A](#), this internal field, coupled with interfacial band bending and electrostatic potential differences, synergistically establishes the foundation for directional charge migration. Mechanistically, an S-scheme system couples an OP with a RP. Upon illumination, the synergistic driving forces compel photogenerated carriers to follow an S-shaped transport route, where low-energy electrons in the conduction band of the OP preferentially recombine with low-energy holes in the valence band of the RP. Consequently, high-energy electrons with robust reducing capabilities (in the RP) and high-energy holes with potent oxidizing capabilities (in the OP) are selectively preserved and directed to their respective active sites. This selective retention mechanism inherently achieves both spatial charge separation and directional charge transfer. Building upon this concept, dual S-scheme heterojunctions, as illustrated in [Figure 6B](#), further introduce multi-interface coupling to construct cascade built-in electric fields, enabling stepwise directional charge transfer across different interfaces. In such systems, the photogenerated carriers undergo hierarchical recombination of low-energy electrons and holes at multiple junctions, while high-energy electrons and holes are progressively separated and enriched at spatially distinct active sites. This multi-level charge regulation not only enhances charge separation efficiency but also maximizes the preservation of strong redox potentials, thereby offering superior performance in driving multi-electron CO₂ reduction reactions. Ultimately, the defining thermodynamic advantage of the S-scheme heterojunction is its ability to fully preserve the original, highly active band edge positions of the constituent semiconductors. This circumvents the energy level compromises typical of conventional heterojunctions, ensuring that the composite maintains sufficient thermodynamic potential to simultaneously drive both CO₂ reduction and complementary oxidation half-reactions.

Recently, Xu *et al.*^[63] reported S-scheme CeO₂/Bi₂S₃ photocatalysis with palladium-catalyzed carbonylation. The CeO₂/Bi₂S₃ system demonstrated a good yield of CO₂ into CO and a high selectivity of approximately 98%, which highlights the comprehensive advantages of the S-scheme structure in enhancing charge separation and product utilization value. Similarly, Zhang *et al.*^[64] engineered an interfacial S-scheme heterojunction by anchoring ultrasmall copper phosphosulfide onto 2D g-C₃N₄. Fourier transform infrared (FT-IR) spectroscopy confirmed the formation of interfacial P-N chemical bonds, which act as highly efficient charge transport channels. This synergistic coupling of chemical bonding and S-scheme transport elevated the CO generation rate to eight times that of pristine g-C₃N₄.

Beyond enhancing overall activity, S-scheme heterojunctions demonstrate a remarkable capability to regulate product selectivity in photocatalytic CO₂ reduction. As summarized in [Figure 6C](#), unlike conventional systems where product distribution is primarily dictated by intrinsic material properties, S-scheme architectures fundamentally reshape the reaction kinetics of competing reduction routes.

Because the mechanism selectively preserves high-energy electrons with highly negative potentials, it provides the robust thermodynamic driving force required for complex, multi-electron transfer processes. This ensures a continuous supply of active electrons to drive the formation of deeply reduced species, such as CH₄, CH₃OH, and C₂+ hydrocarbons, rather than being limited to two-electron products. Consequently, by finely tuning the band alignment, interfacial electric fields, and catalytic active sites, researchers can effectively modulate the competitive pathways of CO₂ reduction.

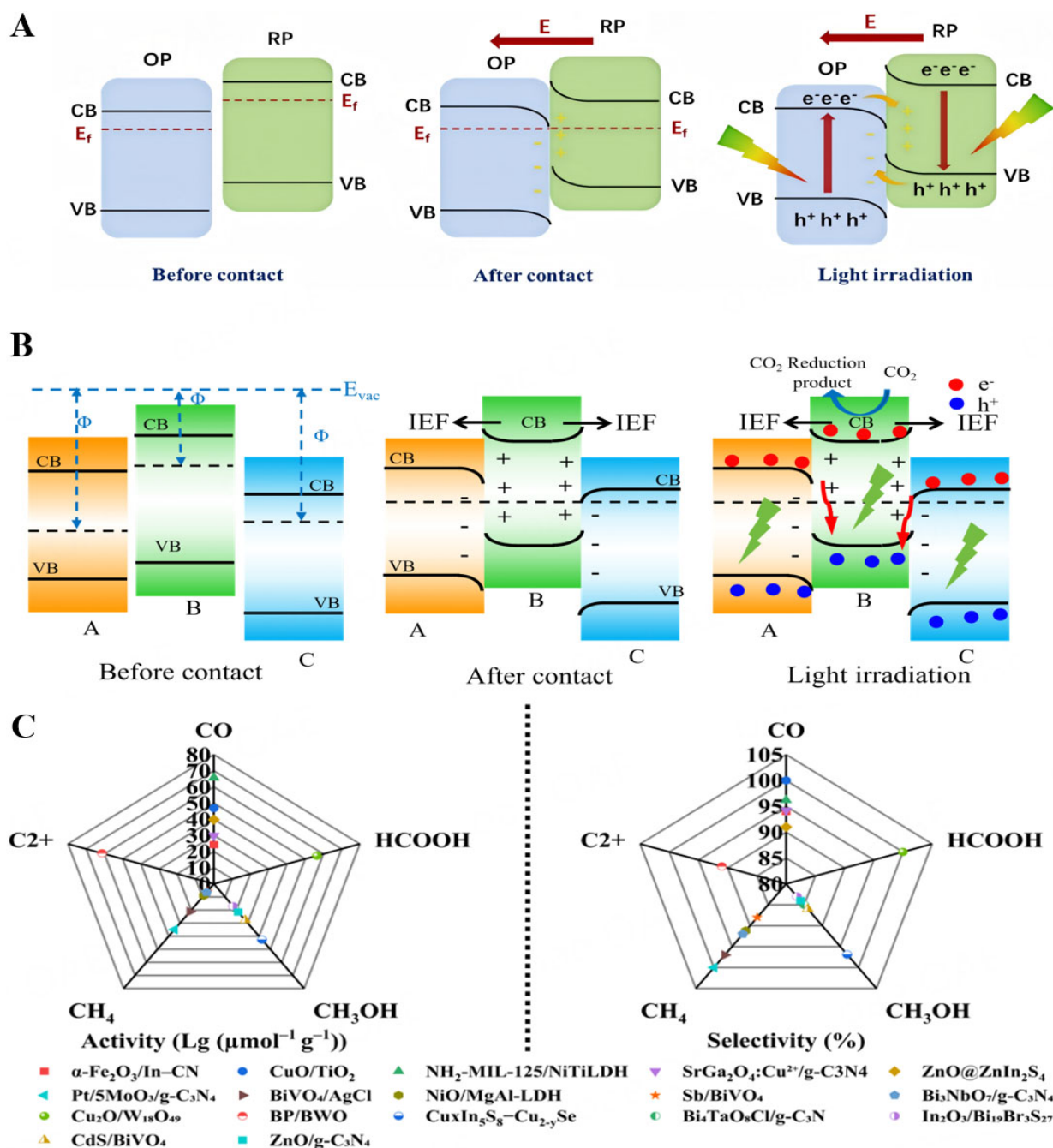


Figure 6. Schematic illustration of the S-scheme heterojunction mechanism: (A) Schematic illustration of the formation mechanism of an S-scheme heterojunction. Reprinted from Ref.^[54] under the CC BY 4.0 license; (B) Schematic illustration of dual S-scheme charge carrier transfer process; (C) Radar plot comparing catalytic activity (expressed as the logarithm of production rate per mass of catalyst) and selectivity among carbon-containing products. CB: Conduction band; VB: valence band; IEF: internal electric field.

The product distributions highlighted in Figure 6C provide a comparative overview of the photocatalytic CO_2 reduction performance of representative S-scheme heterojunctions. The catalytic activity and product selectivity parameters were collected from representative studies reported in Refs.^[65–81] and further extracted, reorganized, and visualized by the authors to facilitate cross-comparison among different material systems. The comparison indicates that S-scheme heterojunctions exhibit diverse hydrocarbon production capabilities, which are closely associated with the synergistic regulation of charge separation, redox potential retention, and surface reaction kinetics. By promoting the accumulation of high-energy electrons and

stabilizing key intermediates (e.g., $^*\text{CO}$ and $^*\text{CHO}$), S-scheme heterojunctions can facilitate subsequent hydrogenation and C-C coupling processes. These characteristics highlight the potential of S-scheme architectures as versatile platforms for the selective synthesis of value-added solar fuels.

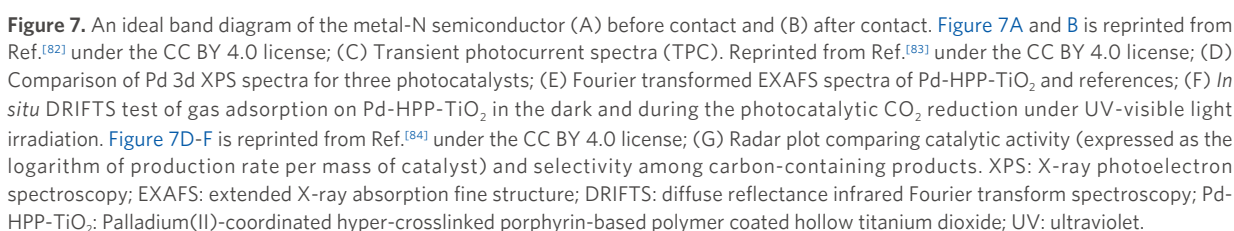
Schottky heterojunction

Schottky heterojunctions are highly effective in enhancing charge separation for photocatalytic CO_2 reduction (PCR). For instance, when a metal contacts an n-type semiconductor [Figure 7A and B]^[82], the metal's typically higher work function drives spontaneous electron transfer across the interface. To satisfy thermodynamic stability, electrons flow from the semiconductor to the metal until their Fermi levels equilibrate. This interfacial equilibration induces local band bending and establishes a built-in electric field, simultaneously forming a depletion layer within the semiconductor. Consequently, a Schottky barrier emerges at the interface, defined by $\Phi_B = \Phi_M - X_{\text{SM}}$, where Φ_M represents the work function of the metal and X_{SM} is the electron affinity of the semiconductor. This barrier fundamentally rectifies charge transport by inhibiting the reverse flow of electrons from the metal back to the semiconductor. As a result, the metal acts as a robust "electron sink", facilitating targeted electron accumulation on its surface and effectively suppressing deleterious charge recombination.

These enhanced photocatalytic performances can be attributed to the favorable interfacial electronic interactions in Schottky structures, which facilitate electron transfer and promote surface reaction processes during CO_2 reduction. Crucially, the functional utility of Schottky junctions extends beyond mere charge separation. In PCR systems, the metal component frequently serves as the primary catalytic active site and CO_2 adsorption-activation center, dynamically regulating the adsorption configuration and electronic structure of key intermediates. Furthermore, when plasmonic metals (e.g., Au, Ag) are employed, localized surface plasmon resonance (LSPR) dramatically amplifies optical absorption. This synergistic coupling of hot-electron injection and the Schottky barrier effect fundamentally fortifies interfacial charge separation and modulates product selectivity. These kinetic advantages are thoroughly corroborated by electrochemical characterizations. As illustrated in Figure 7C, the enhanced transient photocurrent response and diminished electrochemical impedance radius collectively signify improved charge separation efficiency and accelerated interfacial transfer dynamics, ultimately elevating overall PCR performance.

Beyond charge transfer kinetics, nuanced interfacial electronic structures and corresponding reaction mechanisms can be elucidated via advanced *in situ* spectroscopy [Figure 7D-F]. X-ray photoelectron spectroscopy (XPS) and extended X-ray absorption fine structure (EXAFS) mapping provide deep insights into the electronic interactions and coordination environments of the metal sites. Complementarily, *in situ* diffuse reflectance infrared Fourier transform spectroscopy (DRIFTS) visualizes the dynamic adsorption and photo-activation of CO_2 molecules^[84]. Together, these results substantiate that Schottky metal sites operate dually as efficient electron sinks and active centers for intermediate stabilization.

To further validate the general advantages of Schottky architectures beyond individual case studies, a comprehensive literature comparison was performed [Figure 7G]. The integrated radar plot summarizes the reported photocatalytic CO_2 reduction performance of representative Schottky systems. The integrated radar plot in Figure 7G provides a comparative overview of the reported photocatalytic CO_2 reduction performance of representative Schottky architectures. The catalytic activity and product selectivity parameters were collected from representative studies reported in Refs.^[51,83-100] and subsequently extracted, reorganized, and visualized by the authors to enable comparison among different Schottky-type systems and their corresponding single-component counterparts. The comparison indicates that Schottky architectures frequently exhibit enhanced reported catalytic performance, particularly in terms of carbon-containing



Guided by these mechanistic insights, coupling single atoms or highly dispersed nanometals with semiconductors has emerged as a premier strategy for optimizing charge dynamics. For example, in an Ag single-atom anchored TiO₂ system^[101], macroscopic electron trapping-induced coloration is harnessed to

Table 1. Comparison of heterojunction architectures for charge separation in photocatalytic CO₂ reduction

Type	Charge separation mechanism	Carrier transfer pathway	Redox capability	Key advantage	Limitation
Z-scheme	Selective recombination of low-energy carriers	High-energy electrons and holes retained in respective bands	Strong	Maintains strong redox ability and improves charge utilization	Requires precise interface design
S-scheme	Built-in electric field-driven selective recombination	Directional recombination via band bending and Fermi level equilibration	Strong	Combines efficient separation with high redox potential	Mechanism identification can be complex
Schottky	Electron extraction via metal-semiconductor interface	Electrons transferred to metal and trapped by Schottky barrier	Strong	Promotes directional electron transfer and suppresses surface recombination	Limited hole utilization and depends on metal properties

bolster electron storage and migration. The isolated Ag atoms act as dominant active sites to stabilize C₁ intermediates, while adjacent Ti sites facilitate H₂O activation for proton supply, culminating in highly selective CH₄ generation^[101]. Molecular and polymeric photocatalytic platforms analogously benefit from Schottky junction integration to direct charge flow and maximize surface utilization. In benchmark Ag/g-C₃N₄ systems^[102], loading high-work-function metals onto the polymer surface constructs stable Schottky interfaces. This architecture prompts the preferential injection and surface accumulation of photogenerated electrons at the metal sites, severely curtailing bulk-phase recombination and prolonging carrier lifetimes, thereby enhancing the CO production rate significantly beyond that of pristine g-C₃N₄. Moreover, integrating 2D conductive phases, such as MXenes, introduces rapid electron transport channels that augment the interfacial electric field while minimizing contact resistance. For instance, the ZnCr-LDH/Ti₃C₂T_x Schottky system^[85], exhibits exceptional CO yield and selectivity under simulated sunlight, validating the dual role of 2D conductive phases in improving electron extraction and reinforcing the Schottky barrier. Additionally, coupling Schottky junctions with defect engineering further amplifies these kinetic efficiencies. In ultrathin SnS₂ layers enriched with sulfur vacancies, the defect sites dramatically expedite proton and electron supply for the complementary water oxidation half-reaction^[103]. By synergizing the electron utilization efficiency of the Schottky junction with the enhanced oxidation kinetics of the defect sites, this dual-modulation approach bypasses the kinetic bottlenecks associated with half-reaction mismatches, amplifying the CO generation rate to nearly eight times that of the pristine material. Consequently, contemporary research on Schottky heterojunctions for PCR prioritizes three main trajectories: (i) transitioning from metal nanoparticles to single-atom or atomic-level dispersions to achieve highly efficient electron extraction and precisely defined active sites; (ii) integrating 2D conductive phases as electron transport layers to construct stronger built-in electric fields and minimize interfacial resistance; and (iii) synergizing Schottky architectures with defect engineering to mutually optimize barrier heights, charge separation, and half-reaction kinetic matching. Collectively, these advancements establish Schottky heterojunctions as highly tunable interfacial modules for dictating charge and reaction pathways.

To consolidate this mechanistic landscape, the fundamental characteristics of various heterojunction architectures are summarized in Table 1. The core divergence among these designs lies in their driving forces for charge transfer and the consequent impact on redox potential preservation. Conventional type-II heterojunctions rely on standard band alignment, which ensures spatial separation but inevitably sacrifices critical redox energy. In contrast, both Z-scheme and S-scheme systems employ selective recombination mechanisms that preserve high-energy charge carriers necessary for multi-electron CO₂ reduction. Conversely, Schottky heterojunctions are uniquely driven by work function differentials, establishing robust interfacial barriers that enforce directional electron extraction and maximize charge utilization efficiency.

This progressive evolution from simple band-aligned separation toward advanced, thermodynamics-preserving charge regulation highlights the critical necessity of coupling efficient carrier separation with robust redox potential preservation, aligning perfectly with the overarching charge evolution framework conceptualized in Figure 4.

Defect engineering

Defect engineering focuses on deliberately introducing appropriate lattice defects to precisely regulate the local electronic structure and charge-transfer pathways of semiconductors, thereby achieving highly efficient control over charge capture, storage, and reuse^[104]. These defects can induce localized charge density redistribution, modulate band bending, tune built-in electric fields, and dictate the binding affinity and selectivity for surface-adsorbed species. For example, electron-rich defects can enhance molecular adsorption and activate chemical bond bending, thus promoting the formation of critical intermediates (e.g., $^*\text{CO}_2^-$ and $^*\text{COOH}$). Concurrently, hole-trapping centers can suppress non-radiative recombination, ensuring a higher fraction of photogenerated carriers participate in surface catalytic reactions^[105]. Furthermore, the localized polarization and electron-rich microenvironments induced by these defects significantly facilitate direct electron transfer to CO_2 molecules. This establishes a highly favorable thermodynamic landscape for subsequent reduction steps, including C-O bond cleavage and C-H bond formation^[106]. By systematically regulating defect concentration, type, trap depth, and spatial distribution, researchers can precisely optimize charge capture efficiency, migration pathways, and interfacial reaction kinetics—all without perturbing the intrinsic band edges of the host photocatalyst^[107]. Notably, when defect engineering is synergized with complementary strategies, such as interfacial electric field modulation or heterojunction construction, a powerful dual-mechanism emerges: defects serve as localized hubs for electron capture and redistribution, while the enhanced built-in electric field drives rapid, directional charge transfer. Together, these synergistic effects fundamentally elevate overall charge separation and utilization efficiencies in photocatalytic systems.

Vacancy defects

Vacancy defects are one of the most common intrinsic defects in crystalline materials. They refer to the defect formed when lattice nodes lack atoms^[108]. Among them, oxygen vacancies (V_{O}), sulfur vacancies (V_{S}), and metal vacancies (V_{M}) are the most representative. In metal oxides (e.g., TiO_2 , CeO_2 , In_2O_3 , ZnO), V_{O} are typically generated via the detachment of lattice oxygen atoms, leaving behind undercoordinated metal centers and localized electron-rich microenvironments^[109, 110]. These localized states serve a dual function: they act as electron-trapping centers that prolong carrier lifetimes by suppressing non-radiative recombination, while simultaneously serving as prime active sites for the adsorption and activation of CO_2 and critical intermediates (e.g., $^*\text{CO}_2^-$, $^*\text{COOH}$)^[111]. Advancing beyond empirical observations, recent *in situ* and theoretical studies provide a deeper mechanistic understanding. For example, V_{O} in CeO_2 induce the formation of active Ce^{3+} species, which lower the energy barrier for the initial single-electron injection into CO_2 via localized charge redistribution^[112]. Rather than acting merely as static binding sites, these vacancies function as dynamic electron reservoirs that dictate the spatiotemporal distribution of interfacial electrons. However, excessive vacancy concentrations can introduce deep trap states that inadvertently accelerate charge recombination. Consequently, contemporary research has pivoted from indiscriminately maximizing vacancy density toward the precise tuning of vacancy type, trap depth, and spatial distribution to achieve an optimal synergy between charge dynamics and catalytic activity.

Similarly, sulfur vacancies (V_{S}) are predominantly engineered in 2D transition metal dichalcogenides and sulfides (e.g., MoS_2 , WS_2 , and SnS_2). The introduction of V_{S} breaks local crystalline symmetry, exposing previously inert metal centers and substantially amplifying surface reactivity without perturbing the macroscopic band structure^[113]. Recent studies reveal a dual regulatory role for V_{S} ^[114]: they introduce shallow defect states that facilitate electron trapping to extend carrier lifetimes, and they trigger an electronic

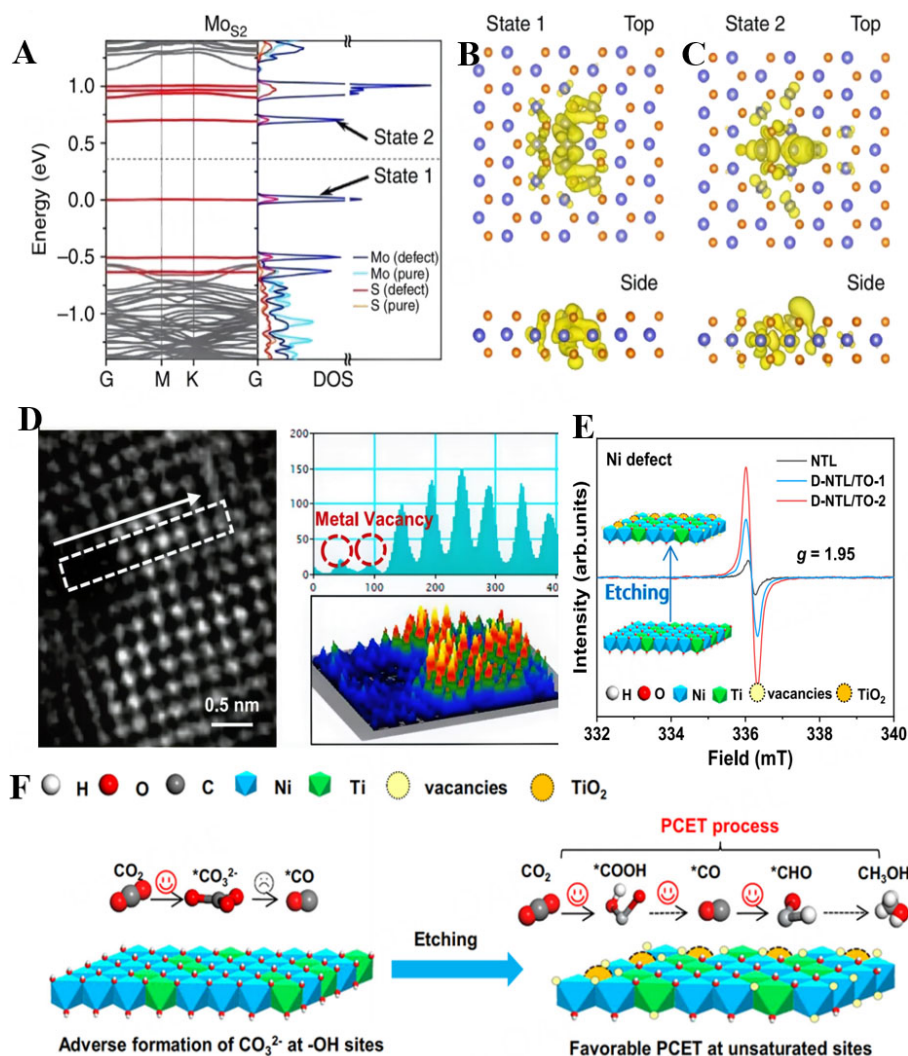


Figure 8. Evidence for vacancy defects in modulating electronic structure and CO₂ reduction pathways: (A) Band structure and corresponding density of states (DOS) of antisite defect MoS₂. The grey bands are from normal lattice sites, similar to conduction band and valence band of perfect monolayer, while the discrete red bands show the localized defects states. The DOS is projected onto the atoms around the defect (defect) and those in the middle plane of two adjacent defects (pure), respectively. The grey dashed line indicates the position of the Fermi Level. (B, C) Real-space distribution of the wave functions of the two defect states below and above the Fermi energy. Figure 8A-C is reprinted from Ref.^[113] under the CC BY 4.0 license; (D) AC-TEM image with corresponding 3D surface plot and atomic profile of white-dotted area, and m EDS-STEM element mapping of D-NTL/TO-2. (E) EPR spectra of bulk NTL and acid-etched derivatives. (F) Schematic illustration of the CO₂ activation mechanism on NTL and D-NTL/TO. Figure 8D-F is reprinted from Ref.^[115] under the CC BY-NC-ND license. D-NTL/TO: Defective NiTi-layered double hydroxide/titanium dioxide; PCET: proton-coupled electron transfer; AC-TEM: Aberration-corrected transmission electron microscopy; EDS-STEM: energy-dispersive X-ray spectroscopy—scanning transmission electron microscopy; EPR: electron paramagnetic resonance.

reconstruction of adjacent metal atoms, which strengthens CO₂ binding affinity and expedites electron injection. Density of states (DOS) analyses provide direct evidence for this V_s-mediated electronic reconstruction. As depicted in Figure 8A-C, the introduction of V_s generates distinct mid-gap localized states (denoted as State 1 and State 2) originating from the rehybridization of metal d and chalcogen p orbitals. These electron-rich mid-gap states effectively accumulate photogenerated electrons, mitigating rapid recombination and providing an energetically favorable pathway for charge transfer to adsorbed CO₂^[113]. Catalytically, these electron-enriched V_s sites promote the back-donation of electrons into the antibonding orbitals of CO₂, accelerating the formation of key intermediates such as *CO₂⁻ and *COOH. This underscores that sulfur vacancies actively orchestrate interfacial charge transfer and molecular activation rather than functioning as passive electron traps. Consequently, V_s engineering is now a cornerstone design strategy for sulfide-based photocatalytic systems.

In contrast to anion vacancies, metal vacancies (V_M) generally demand higher formation energies, typically requiring harsh annealing or reducing atmospheres, but exert a far more profound impact on the electronic structure and catalytic behavior of the host material^[104]. V_M disrupt the local coordination environment, inducing substantial charge redistribution and lattice polarization that enhance both carrier transport and interfacial kinetics. As illustrated in Figure 8D, high-resolution intensity mapping reveals that structurally unsaturated V_M sites trigger pronounced local lattice distortions. These coordination-unsaturated centers induce spatially confined electronic perturbations, fostering charge localization under reaction conditions. This electronic modulation is further corroborated by electron paramagnetic resonance (EPR) spectroscopy [Figure 8E], where a distinct signal at $g = 1.95$ directly evidences the formation of paramagnetic defect centers associated with Ni vacancies. The intensity evolution of this signal across different treatment conditions highlights that V_M concentration and electronic activity can be dynamically manipulated via post-etching strategies. Catalytically, these V_M -rich domains critically govern adsorption energetics and intermediate evolution. As summarized in Figure 8F, a defective NiTi-TiO₂ system effectively suppresses the formation of thermodynamically stable yet catalytically inert carbonate species, selectively stabilizing the *COOH intermediate, a widely recognized rate-determining precursor for CO₂-to-methanol conversion^[115]. This pathway reconfiguration stems from the synergy between V_M -induced electronic state redistribution (enhancing electron availability) and localized geometric distortions (optimizing proton accessibility). Crucially, this modulation transcends static adsorption to actively accelerate PCET kinetics. By shortening the effective distance between charge carriers and intermediates, these unsaturated sites lower the activation barriers for sequential hydrogenation steps (*CO₂ → *COOH → *CO → *CHO). Furthermore, from a carrier dynamics perspective, V_M in metal oxides (e.g., WO₃, TiO₂, Bi₂O₃) decrease the effective mass of holes, thereby boosting their mobility and expediting bulk-to-surface carrier migration^[116]. In transition metal oxides like Co₃O₄, V_M promote the formation of high-valence metal species, further driving oxidation kinetics^[117]. Thus, metal vacancies operate as multifunctional catalytic motifs that concurrently optimize charge separation, intermediate stabilization, and reaction pathway selectivity.

Ultimately, the formation and behavior of these vacancy defects are intimately governed by the interplay of material composition, crystal structure, and synthetic conditions (e.g., temperature and atmospheric composition). Thermodynamically, vacancy generation must overcome binding-energy-dependent barriers; kinetically, it is dictated by atomic diffusion during thermal or reductive treatments. Table 2 comprehensively summarizes recent advances in vacancy-engineered photocatalysts for PCR. These benchmarks collectively illustrate that whether leveraging V_O in metal oxides, V_S in sulfides, or more complex V_M and dual-defect architectures ($V_M + V_S$), vacancy engineering remains a paramount strategy. It not only suppresses deleterious charge recombination but also fundamentally dictates the surface reaction pathways and product selectivity in photocatalytic CO₂ reduction.

Doping defects

Doping defects, introduced via the interstitial or substitutional incorporation of metallic or non-metallic heteroatoms, profoundly reconfigure the electronic band structure, local coordination environments, and charge dynamics of host lattices^[109]. Unlike intrinsic vacancies, doping provides a highly tunable platform to actively dictate the spatiotemporal behavior of photogenerated carriers and optimize interfacial catalytic performance.

Non-metal dopants (e.g., N, C, S, P, B) primarily modulate the valence band architecture, inducing electronic state reconstruction and generating localized impurity states^[106]. Beyond merely broadening the optical absorption spectrum, non-metal doping fundamentally governs interfacial charge transfer. For example, synergistic P-S co-doping transforms polymeric g-C₃N₄ into a narrow-bandgap “black” carbon nitride. This architectural modification not only extends light harvesting into the near-infrared regime but also instigates

Table 2. Summary of vacancy-defect-engineered photocatalysts for CO₂ reduction

Material	Vacancy type	Main product	Light sources	Sacrificial agents	AQE (%)	Rate $\mu\text{mol}\cdot\text{g}^{-1}\cdot\text{h}^{-1}$	Ref.
TiO _{2-x}	V _O	CH ₄	300 W Xe lamp	H ₂ O	0.12 (365 nm)	41.8	[118]
Pd/CeO ₂	V _O	CO	300 W Xe lamp	H ₂ O	-	210.9	[119]
Bi ₁₂ O ₁₇ Cl ₂	V _O	CO	300 W Xe lamp	H ₂ O	-	64.3	[120]
BiOCl	V _O	CO	300 W Xe lamp	H ₂ O	-	14.51	[121]
BiOBr	V _O	CO	300 W Xe lamp $\lambda \geq 420$ nm	H ₂ O	-	122.38	[122]
BWO	V _O	CO	300 W Xe lamp	H ₂ O	0.00394 (500 nm)	18.73	[123]
MoO _{2-x}	V _O	CO	250 W high-pressure mercury lamp	H ₂ O	-	62.75	[124]
Bi/BiOBr	V _O	CO	200 mW/cm ² Xe lamp	H ₂ O	-	251.2	[125]
Bi ₂ MoO ₆ @In ₂ S ₃	V _O	CO	300 W Xe lamp $\lambda \geq 420$ nm	H ₂ O	2.11 (420 nm)	28.54	[126]
ReS ₂ /CdS	V _S	CO	300 W Xe lamp $\lambda \geq 420$ nm	H ₂ O	-	7.1	[127]
ZIS	V _S	CO	300 W Xe lamp $\lambda \geq 420$ nm	TEOA	-	61.94	[128]
ZnS/OMNC	V _S	CO	300 W Xe lamp $\lambda \geq 420$ nm	H ₂ O	-	712.1	[129]
SnS ₂	V _S	CO	300 W Xe lamp	H ₂ O	-	25.71	[103]
Au/CdS-SV	V _S	CO	300 W Xe lamp $\lambda \geq 400$ nm	TEOA	-	12.48	[130]
ZnIn ₂ S ₄	V _M	CO	300 W Xe lamp $\lambda \geq 420$ nm	H ₂ O	2.29 (420nm)	5630	[131]
BiOBr-1	V _M	CO	300 W Xe lamp	H ₂ O	-	71.23	[132]
VBi-BiOBr	V _M	CO	300 W Xe lamp	H ₂ O	-	20.1	[133]
Cu _{1.95} S _{1-x}	V _M +V _S	CH ₄	300 W Xe lamp	H ₂ O	-	12.42	[134]

AQE: Apparent quantum efficiency; ZIS: ZnIn₂S₄; OMNC: ordered mesoporous N-doped carbon; TEOA: triethanolamine.

a redistribution of electron density and local structural polarization, dramatically accelerating CO₂-to-CO conversion by mitigating electron-hole recombination^[135].

Conversely, metal ion doping predominantly engineers shallow energy levels near the conduction band and modulates the Fermi level, steering the directional migration of photogenerated carriers. A quintessential example of metal-doping synergy is the Bi-incorporated In₂O_{3-x}(OH)_y system. Here, the strategic construction of surface frustrated Lewis pairs (FLPs, cooperatively regulated by hydroxyl groups and oxygen vacancies) is further optimized by Bi³⁺ substitution at In³⁺ sites. This targeted metal incorporation finely tunes the surface active motifs, dictating both the thermodynamic activation of CO₂ and the ultimate product selectivity^[136,137].

Extending these principles, dual-doping strategies (e.g., metal/non-metal combinations or dopant-vacancy couplings) elicit robust synergistic effects. For instance, P-K co-doping in g-C₃N₄ amplifies charge separation while driving the multielectron reduction of CO₂ toward CH₄^[138], whereas N-metal co-doping in TiO₂ concurrently broadens spectral responsiveness and enhances product selectivity^[139]. Table 3 systematically catalogs recent benchmarks in doping-engineered photocatalysts, highlighting the macroscopic performance gains achieved through these elemental substitutions.

Table 3. Summary of doped photocatalysts for CO₂ reduction

Material	Vacancy type	Main product	Light sources	Sacrificial agents	AQE (%)	Rate $\mu\text{mol}\cdot\text{g}^{-1}\cdot\text{h}^{-1}$	Ref.
TiO ₂	N	CO	300 W Xe lamp	H ₂ O	-	41.1	[140]
TiO ₂	N+Eu	CH ₄	300 W Xe lamp $\lambda \geq 400$ nm	H ₂ O	9.1 (350 nm)	13.48	[141]
BiOBr	P	CO	300 W Xe lamp	H ₂ O	-	9.13	[142]
BiOBr	N	CO	300 W Xe lamp	H ₂ O	-	18.28	[143]
BiOBr	B	CO	300 W Xe lamp	H ₂ O	-	21.72	[144]
BiOBr	Co	CO	300 W Xe lamp $\lambda \geq 420$ nm	H ₂ O	-	11.71	[145]
BiOBr	Eu	CO	300 W Xe lamp	H ₂ O	-	22.3	[146]
BiOBr-2	Br	CO	300 W Xe lamp	H ₂ O	-	95.6	[147]
ZnIn ₂ S ₄	O	CO	300 W Xe lamp $\lambda \geq 420$ nm	TEOA	-	1680	[148]
ZnIn ₂ S ₄	Er	CH ₄	300 W Xe lamp $\lambda \geq 420$ nm	TEOA	-	6.68	[149]
ZnIn ₂ S ₄	Y	CO	300 W Xe lamp $\lambda \geq 420$ nm	TEOA	0.97 (350 nm)	297.46	[150]
ZnIn ₂ S ₄	Yb	CO	300 W Xe lamp $\lambda \geq 420$ nm	TEOA	-	4.59	[151]
CN	P+F	CO	300 W Xe lamp $350 \text{ nm} \leq \lambda \leq 780 \text{ nm}$	H ₂ O	-	39.9	[152]
CN	P	CO	300 W Xe lamp	H ₂ O	-	31.22	[153]
Au/CN	B+K	CO	300 W UV lamp $\lambda = 400$ nm	H ₂ O	-	11.56	[154]
g-C ₃ N ₄	Eu	CH ₄	300 W Xe lamp	H ₂ O	1.6 (420 nm)	22.8	[155]
CN	Cu+P	CO	300 W Xe lamp	H ₂ O	0.005 (375 nm)	6.01	[156]
Bi ₃ O ₄ Br	Co	CO	300 W Xe lamp	H ₂ O	-	107.1	[157]

AQE: Apparent quantum efficiency; TEOA: triethanolamine.

To bridge these macroscopic catalytic enhancements with atomic-scale charge dynamics, interstitial carbon-doped SnS₂ (SnS₂-C) serves as an elegant mechanistic model [Figure 9A-C]. Interstitial carbon incorporation disrupts local structural symmetry, inducing pronounced lattice microstrain and charge polarization. High-resolution XPS [Figure 9A and B] unveils a distinct positive shift in the binding energies of both Sn 3d and S 2p core levels. This phenomenon signifies partial electron depletion from the host lattice, indicating that interstitial carbon acts as a potent electron-withdrawing locus. Such localized electron redistribution generates a robust built-in electric field that drives the spatial separation of bulk electron-hole pairs, funneling them directionally toward the catalytic interface. Furthermore, as depicted in Figure 9C, the synergy between carbon-induced lattice strain and electronic modulation broadens visible-light harvesting while establishing energetically favorable pathways for charge injection into adsorbed CO₂. Ultimately, this paradigm underscores that doping defect engineering transcends simple bandgap narrowing; it acts as a comprehensive tool to construct internal electric fields, optimize charge routing, and fundamentally dictate the fate of photogenerated carriers at the atomic level.

Cocatalyst modification

Loading cocatalysts on the photocatalyst surface can enhance charge separation, improve charge transfer, or adjust the reaction pathway, thereby improving the overall PCR performance. Generally, cocatalysts are classified into metal cocatalysts and non-metal cocatalysts^[109].

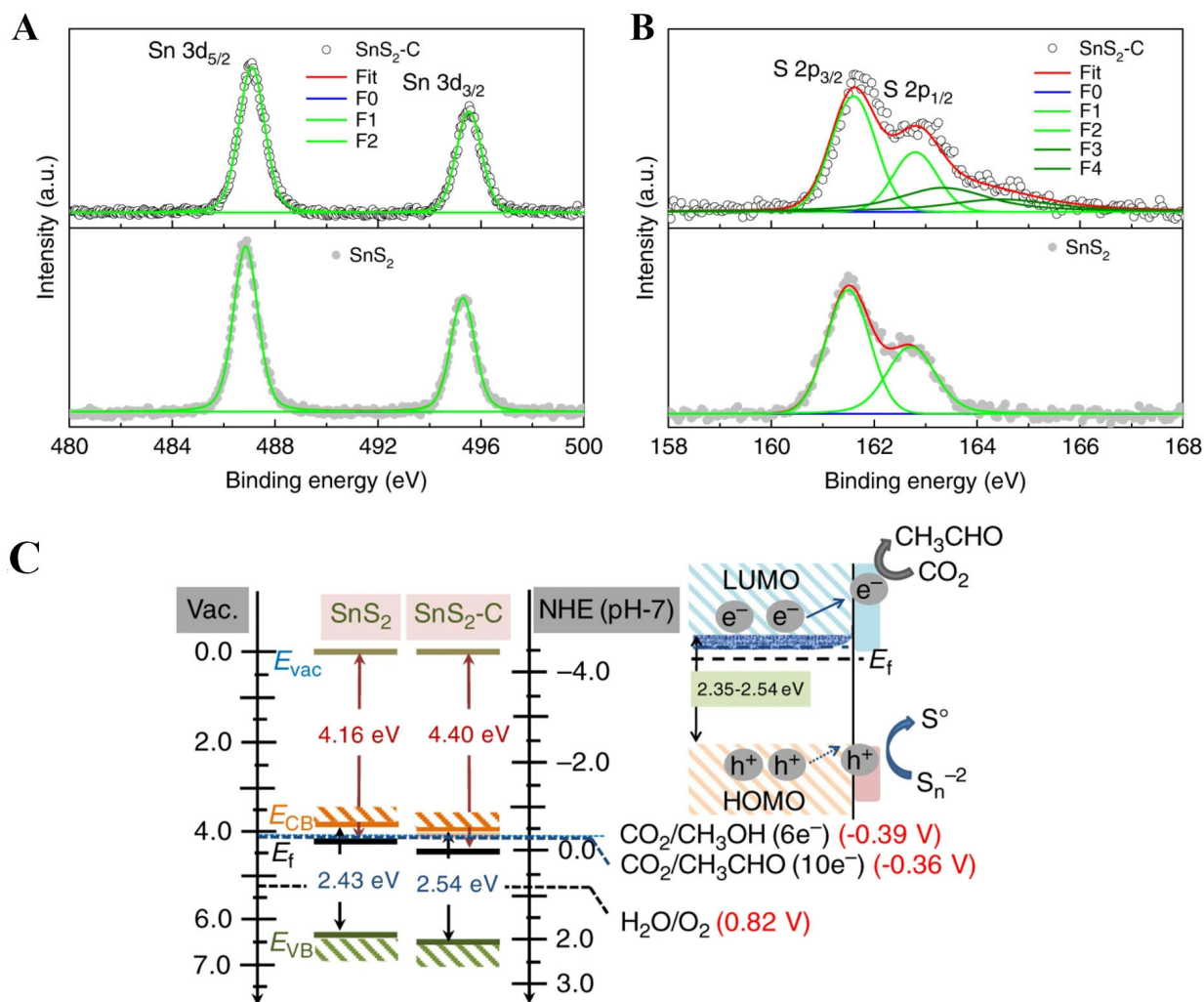


Figure 9. (A) High-resolution XPS Sn 3d spectra of SnS₂-C and SnS₂; (B) High-resolution XPS S 2p spectra of SnS₂-C and SnS₂; (C) Band edge positions and photocatalytic reaction mechanism: Comparative band diagram of SnS₂-C and SnS₂, together with a proposed electron-hole separation of photo-excited electron-hole pairs in SnS₂-C. Figure 9A-C is reprinted from Ref.^[158] under the CC BY 4.0 license. SnS₂-C: Carbon-doped SnS₂; CB: conduction band; VB: valence band; NHE: normal hydrogen electrode; LUMO: lowest unoccupied molecular orbital; HOMO: highest occupied molecular orbital; XPS: X-ray photoelectron spectroscopy.

Cocatalysts serve as indispensable structural motifs in photocatalytic systems, bridging the spatial gap between bulk charge separation and interfacial catalytic turnover. Among these, metal nanostructures remain a cornerstone, functioning as potent electron sinks. High-work-function noble metals (e.g., Pt, Au, Pd) equilibrate with the Fermi level of n-type semiconductors upon contact^[159], inducing spontaneous band bending and the formation of a robust Schottky junction [Figure 10A]. This built-in potential barrier actively circumvents electron backflow, funneling photogenerated electrons directionally toward the metallic surface where CO₂ reduction transpires. Furthermore, plasmonic cocatalysts (e.g., Ag, Au) harness localized surface plasmon resonance (LSPR) to inject hot electrons directly into the semiconductor conduction band, synergistically amplifying the overall quantum efficiency alongside Schottky-directed charge extraction^[102,159].

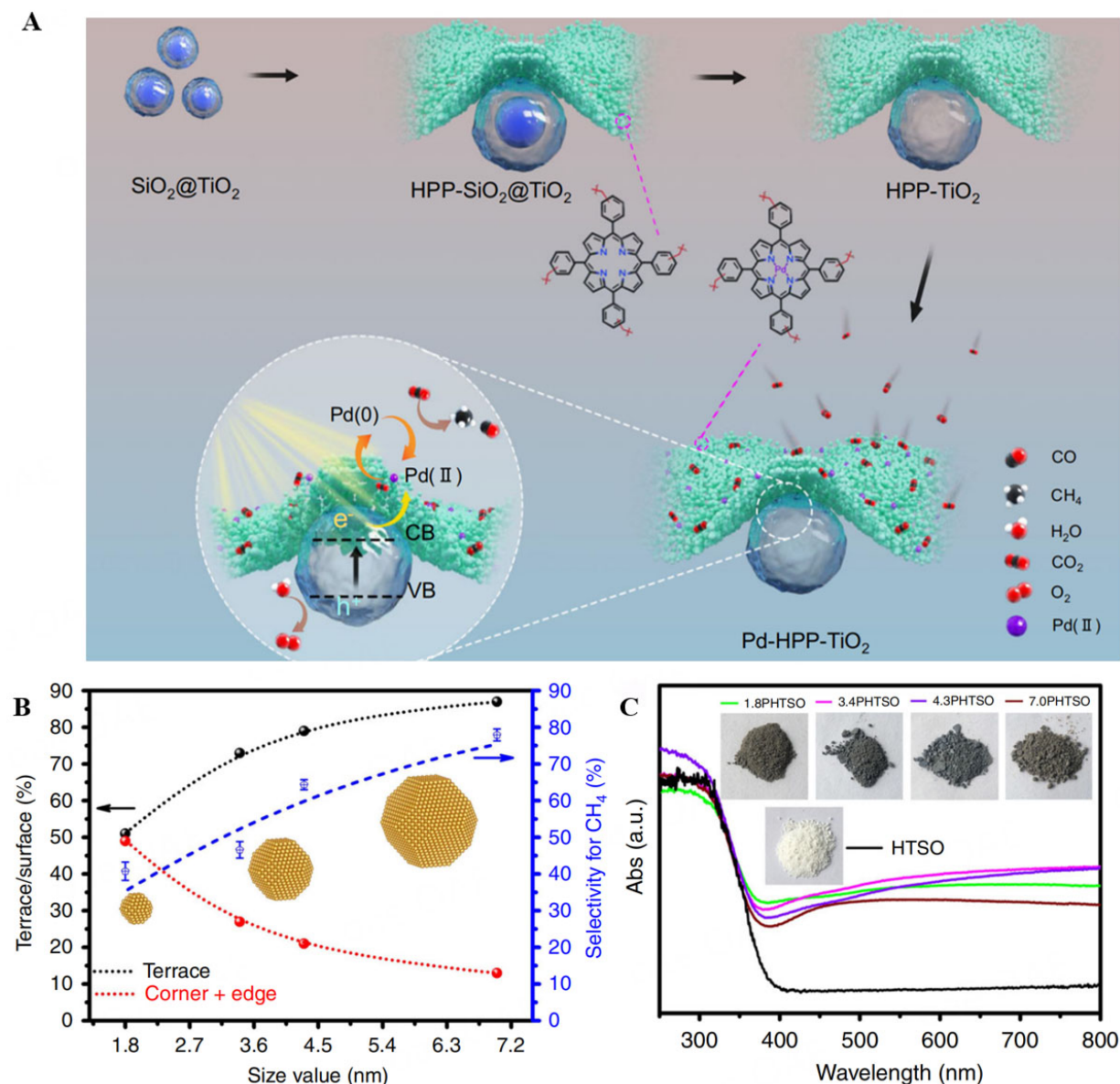


Figure 10. (A) Synthesis of porous Pd-HPP-TiO₂ and the possible mechanism of photocatalytic CO₂ reduction. Reprinted from Ref.^[64] under the CC BY 4.0 license. (B) Correlations between the selectivity for CH₄ and surface site proportion as functions of the size of Pt NPs in xPHTSO (x = 1.8, 3.4, 4.3, and 7.0). (C) UV-Vis DRS spectra and the corresponding samples' photos (inset). Figure 10B and C is reprinted from Ref.^[160] under the CC BY 4.0 license. HPP: Hyper-crosslinked porphyrin-based polymer; CB: conduction band; VB: valence band; PHTSO: Pt-loaded hollow TiO₂ with a silica overlayer; NP: nanoparticle; UV-Vis: ultraviolet-visible; DRS: diffuse reflectance spectroscopy.

Crucially, the catalytic efficacy of these metallic sinks is exquisitely sensitive to their geometric dimensions [Figure 10B and C]. Downsizing nanoparticles maximizes the population of low-coordinated surface atoms (corner and edge sites), augmenting electron trapping efficiency^[160]. However, this miniaturization simultaneously alters the binding energetics of crucial intermediates, occasionally compromising target product selectivity. Conversely, larger nanoparticles predominantly expose terrace sites that thermodynamically favor specific multistep pathways, such as deep reduction to CH₄. This intricate trade-off dictates that metal cocatalyst design must transcend mere work-function matching, necessitating precise geometric tailoring to harmonize charge extraction kinetics with surface reaction specificity.

Pushing this geometric limit to the atomic scale, single-atom catalysts (SACs), such as Cu, Ni, and Co, have emerged as a transformative frontier. Unlike their nanoparticulate counterparts, SACs offer maximum atom

utilization, well-defined coordination environments, and the ability to modulate local electronic structures without fundamentally perturbing the host semiconductor's band architecture^[161]. In model systems like Ag-anchored TiO₂, isolated Ag atoms operate as exclusive electron-trapping loci, substantially extending carrier lifetimes while stabilizing critical C₁ intermediates. The adjacent Ti sites synergistically facilitate water oxidation and proton supply, dictating a highly selective CO₂-to-CH₄ conversion pathway^[101]. Similarly, anchoring isolated Cu species on polymeric carbon nitride (Cu-PCN) directly mediates electron transfer to adsorbed CO₂. Theoretical and *in situ* spectroscopic investigations corroborate that these atomic Cu sites significantly lower the activation barrier for *COOH formation, the rate-determining step, thereby steering the reaction with near-unity selectivity toward CO^[151-153].

Complementing metallic systems, non-metallic and quasi-metallic cocatalysts leverage extended π -conjugated networks and high carrier mobilities to orchestrate directional charge routing. Carbon quantum dots (CQDs), acting as dynamic “electron reservoirs”, provide rapid, delocalized electron transport channels that decisively mitigate bulk recombination when integrated into heterostructures (e.g., TiO₂/SrTiO₃ or Z-scheme Bi₁₂O₁₇Cl₂/NiAl-LDH)^[162,163]. Furthermore, covalent organic frameworks (COFs) and conjugated polymers offer unprecedented molecular-level tunability. By constructing continuous interfacial π -electron channels, these materials expedite carrier migration while circumventing the strong, often irreversible metal-intermediate bonds characteristic of inorganic systems^[164]. Consequently, non-metal cocatalysts typically govern intermediate stabilization via soft non-covalent interactions (e.g., π - π or electrostatic forces). By fine-tuning the adsorption energetics of species like *CO, these networks strategically suppress over-hydrogenation pathways, favoring highly selective CO evolution^[165].

Table 4 systematically compiles recent benchmarks in cocatalyst-engineered photocatalysts, detailing the interplay between cocatalyst identity, half-reaction dynamics, and product distribution. Collectively, these studies underscore that rational cocatalyst integration is not merely a strategy to accelerate redox kinetics, but a fundamental prerequisite for bridging spatial charge extraction with precisely tailored interfacial reaction pathways in advanced CO₂ reduction systems.

Microstructure design

Microstructure engineering represents a quintessential strategy in photocatalytic CO₂ reduction (PCR) systems to fundamentally optimize mass diffusion and truncate charge transfer pathways. Dimensionality reduction to the 2D ultrathin limit effectively circumvents bulk recombination by ensuring that structural thicknesses are comparable to, or smaller than, the carrier diffusion length. For instance, atomic-scale Bi₄O₅Br₂ nanosheets exhibit drastically enhanced CO₂-to-CO activity compared to their bulk counterparts, as photogenerated electron-hole pairs are inherently forced to migrate toward surface active sites^[186].

Furthermore, assembling 2D/2D heterojunctions (e.g., LDHs/MoS₂ or MXene/g-C₃N₄) maximizes interfacial contact area while minimizing transport resistance. In such architectures, highly conductive substrates such as MXenes function as robust “electron highways” expediting the directional funneling of photogenerated electrons to surface catalytic loci and markedly suppressing recombination^[187,188].

Transitioning to 3D hierarchical designs, core-shell architectures introduce radial charge partitioning and coaxial transport cascades. By spatially decoupling oxidation and reduction half-reactions across a distinct geometric boundary, these structures foster robust built-in electric fields that orchestrate directional charge flow. Complex configurations, such as TiO₂@NiAl-LDH, Cu₂O@Cu-MOF, and ternary Cu₂O@NiAl-LDH/CQDs systems, vividly demonstrate how rational shell selection concentrates CO₂ molecules at the outermost reactive sites while spatially isolating them from internal charge generation centers, thereby firmly averting reverse charge recombination and enhancing multi-electron C₂ product selectivity^[185,189-191].

Table 4. Summary of cocatalyst-modified photocatalysts for CO₂ reduction

Material	Vacancy type	Main product	Light sources	Sacrificial agents	AQE (%)	Rate $\mu\text{mol}\cdot\text{g}^{-1}\cdot\text{h}^{-1}$	Ref.
Co-Bi ₂ O ₄ Br	single-atom Co	CO	300 W Xe lamp	H ₂ O	-	107.1	[157]
Co-TCPP/Bi ₂ O ₄ Br	molecular Co	CO	300 W Xe lamp	H ₂ O	0.53 (380 nm)	71.3	[166]
CO ₂ N/BiOBr	non-noble nitride	CO	300 W Xe lamp	H ₂ O	-	67.8	[167]
Ni SAs-Bi ₂ O ₄ Br	single-atom Ni	HCOOH	300 W Xe lamp	H ₂ O	-	3632.6	[168]
Cs ₃ Bi ₂ Br ₉ /Co-NG	Co single-atom modified graphene	CO	300 W Xe lamp	H ₂ O	-	123.16	[169]
LaNi-Phen/COF-5	La-Ni bimetallic	CO	500 W Xe lamp	H ₂ O	-	605.8	[170]
Cu-HCOF	atomically dispersed Cu	CO	400-W high-pressure Hg lamp	TEOA	-	960	[171]
AgCo/Al-SrTiO ₃	dual Ag/Co	CO	400-W high-pressure Hg lamp	H ₂ O	0.03 (365 nm)	105.4	[172]
Ag@Cr/Ga ₂ O ₃	Ag-Cr core-shell	CO	300 W Xe lamp	H ₂ O	-	999.2	[173]
Cu/TiO ₂	Cu	CH ₄	300 W Xe lamp	H ₂ O	-	12.52	[174]
PtCuCN	Pt-loaded CuPc/g-C ₃ N ₄	CH ₄	LED light source (365 $\pm$ 10 nm)	TEOA	-	39.8	[51]
Pt-Au/R-TNTs	Pt-Au	CH ₄	UV light irradiation (300 nm < λ < 400 nm)	H ₂ O	17.9 (365 nm)	360	[175]
Pt/TiO ₂	Pt	CH ₄	300 W Xe lamp $\lambda \geq 400$ nm	H ₂ O	-	4.6	[176]
Pd TP/Pd SA-CN	Pd single atoms + twinned Pd nanoparticles	CO	600 mW/cm ²	H ₂ O	-	46.5	[177]
Pd NPs/TiO ₂	Pd	CH ₄	350 W Xe lamp	H ₂ O	-	28.5	[178]
g-C ₃ N ₄ /MoS ₂ /Cu	Cu	CO	300 W Xe lamp	H ₂ O	3.2 (420 nm)	146.7	[179]
Co ₂ In ₂ /CN	Co-In dual single-atom	CH ₄	UV light	H ₂ O	-	18.8	[180]
Pt@CeO ₂ /3D CN	Pt	CO	300 W Xe lamp	TEOA	-	4.69	[181]
Cu _{0.7} Au _{0.3} /TiO ₂	Cu-Au bimetal	CO	300 W Xe lamp	H ₂ O	-	6.08	[182]
Ag-Cu/TiO ₂	Ag-Cu dual	CO	300 W Xe lamp	H ₂ O	-	286.7	[183]
Au ₆ Ag ₄ /TiO ₂	Au-Ag alloy	CO	300 W Xe lamp 320 nm < λ < 780 nm	H ₂ O	0.91 (450 nm)	99.54	[184]
Cu ₂ O@Cu-MOF/TiO ₂	Cu ₂ O/Cu-MOF core-shell	CH ₄	300 W Xe lamp	H ₂ O	-	366	[185]

AQE: Apparent quantum efficiency; TCPP: tetrakis(4-carboxyphenyl)porphyrin; SA: single-atom; TEOA: triethanolamine; R-TNT: Reduced TiO₂ nanotubes; MOF: metal-organic framework; NP: nanoparticle.

Porous and hollow microstructures culminate this spatial optimization by synchronously addressing the “light-harvesting, mass-transfer, and charge-separation” trifecta. As systematically elucidated in Figure 11, the advantages of these geometries are multidimensional. Internally, multi-shelled or hollow cavities induce severe light scattering and multireflection [Figure 11A], drastically amplifying the photon absorption path length. Electronically, the attenuated physical distance from the bulk to the solid-gas interface expedites charge extraction, a phenomenon universally corroborated by quenched steady-state photoluminescence (PL) [Figure 11B] and prolonged carrier lifetimes in time-resolved PL (TRPL) spectra [Figure 11C]. Concurrently, the interconnected porous networks massively expand surface accessibility, translating to superior CO₂ adsorption capacities [Figure 11D]^[194].

These synergistic mechanisms dictate the exceptional performance of practical systems, such as hierarchical porous ZnIn₂S₄ nested heterojunctions, hollow g-C₃N₄@TiO₂ microspheres, and hollow CeO₂ structures^[193–196]. In these configurations, the structural voids not only host enriched active sites but

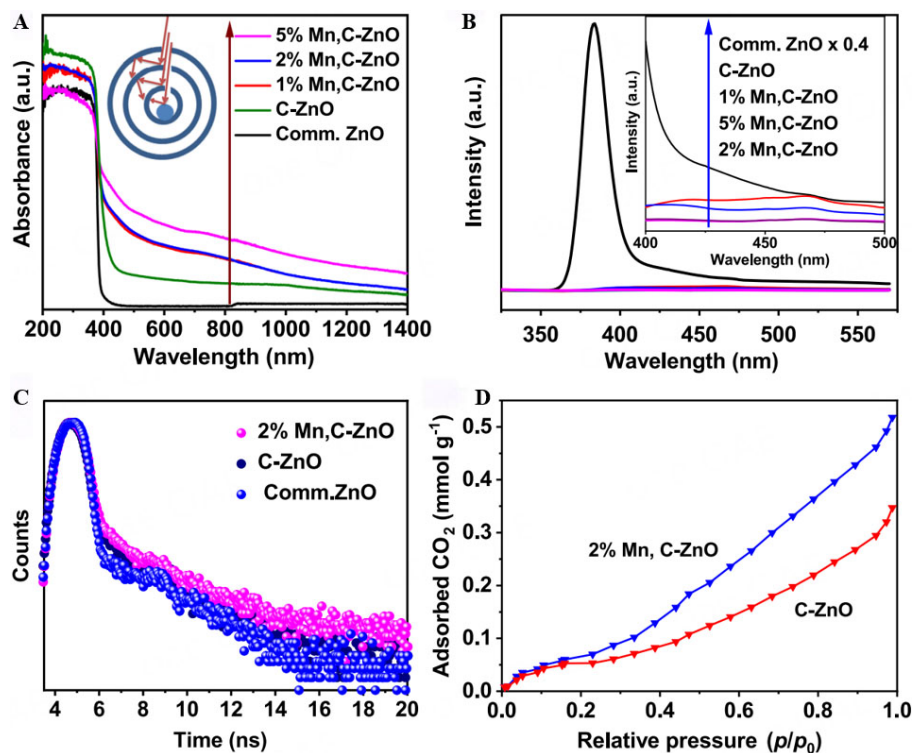


Figure 11. (A) UV-vis light absorption of the prepared photocatalysts and the inset represents the multiple reflection effect of the incident light inside the hollow cavities. (B) PL spectra of photocatalysts at excitation wavelength of 380 nm. (C) TRPL spectra of the samples. (D) CO₂ adsorption isotherm of C-ZnO CTS HSs and 2% Mn, C-ZnO-CTS HS samples. Figure 11 A-D is reprinted from Ref.^[192] under the CC BY 4.0 license. UV-vis: Ultraviolet-visible; PL: photoluminescence; TRPL: time-resolved photoluminescence; CTS HS: core-triple shell hollow sphere.

fundamentally accelerate interfacial reaction kinetics under visible-light irradiation.

Table 5 summarizes state-of-the-art microstructure-engineered photocatalysts. Collectively, these geometric modulations, ranging from atomic 2D exfoliation to hierarchical 3D hollowing, demonstrate that macroscopic catalytic enhancements are inextricably linked to the precise nanoscale management of photon flux, reactant diffusion, and carrier transport vectors.

While the aforementioned strategies ranging from band engineering to microstructural design have individually propelled photocatalytic CO₂ reduction (PCR) forward, a critical appraisal reveals a persistent “single-strategy bottleneck”. For instance, Z- and S-scheme heterojunctions elegantly preserve high-energy charge carriers, yet their efficacy is strictly tethered to interfacial electronic coupling; poorly matched work functions or structural deviations inevitably spawn parasitic recombination pathways. Similarly, while metallic sinks and Schottky junctions expedite charge extraction, their optimization is plagued by a delicate balancing act, as excessive loading invariably triggers light-shielding effects or active-site occlusion. Furthermore, morphological tuning and defect engineering (e.g., oxygen vacancies) effectively shorten diffusion lengths and enhance reactant adsorption, respectively, but they often lack the intrinsic thermodynamic driving force required to independently orchestrate long-distance, directional charge separation.

Consequently, tackling the PCR bottleneck demands a paradigm shift from isolated component optimization to the synergistic integration of multi-dimensional strategies. This synergy transcends a mere mathematical superposition of effects; rather, it orchestrates a holistic reconstruction of charge-transfer networks across

Table 5. Summary of microstructure-engineered photocatalysts for CO₂ reduction

Material	Vacancy type	Main product	Light sources	Sacrificial agents	AQE (%)	Rate $\mu\text{mol}\cdot\text{g}^{-1}\cdot\text{h}^{-1}$	Ref.
Bi ₄ O ₅ Br ₂ -UN	2D nanosheets	CO	300 W Xe lamp	H ₂ O	-	31.57	[186]
TC/g-C ₃ N ₄	2D/2D heterojunction	CO	300 W Xe lamp $\lambda \geq 420$ nm	H ₂ O	-	5.19	[188]
TiO ₂ @CN	core-shell Z-scheme	CO	300 W Xe lamp 320 nm < λ < 780 nm	H ₂ O	1.829 (400 nm)	26.89	[191]
Cu ₂ O@NA-LDH/CQDs	ternary core-shell	C ₂ H ₆	300 W Xe lamp $\lambda \geq 420$ nm	TEOA	-	8.18	[190]
NiAl-LDH@TiO ₂ /Ti ₃ C ₂	hierarchical core-shell	CO	300 W Xe lamp	H ₂ O	0.81 (365 nm)	86.52	[197]
CNGA/CdS	ultrathin porous nanosheets	CO	300 W Xe lamp $\lambda \geq 420$ nm	H ₂ O	0.65 (500 nm)	32.75	[198]
CDs-1/CdS	hierarchical hollow microspheres	CO	300 W Xe lamp	TEOA	-	16.09	[199]
H-Cs ₃ Sb ₂ Br ₉	hollow nanospheres	CO	300 W Xe lamp $\lambda \geq 400$ nm	CH ₃ OH	2.65 (400 nm)	1876.3	[200]
Bi ₄ O ₅ Br ₂	hollow core-shell	CO	300 W Xe lamp	H ₂ O	-	3.61	[201]
NiCO ₂ V ₂ O ₈	hollow nanospheres	CO	300 W Xe lamp $\lambda \geq 420$ nm	TEOA	5.76 (380 nm)	198.65	[202]
g-C ₃ N ₄ /rGO/NiAl-LDHs	layered hierarchical heterojunction	CH ₄	300 W Xe lamp	TEOA	1.36 (380 nm)	20	[203]
BWO	ultrathin nanosheets	CO	300 W Xe lamp	H ₂ O	0.35 (350 nm)	19.45	[204]
F-2D BMO	2D nanosheets	CO	300 W Xe lamp	H ₂ O	1.5 (420 nm)	5.6	[205]
CdS/TiO ₂	hollow microspheres	CH ₄	350 W Xe lamp	H ₂ O	-	27.85	[206]
Cu-HCOF	hollow COF structure	CO	500 W Xe lamp	TEOA	-	960	[171]
Bi-CuO	core-shell	CO	300 W Xe lamp	H ₂ O	-	106.07	[207]
Ag/hollow-TiO ₂	hollow spheres	CH ₄	300 W Xe lamp	H ₂ O	-	6.72	[208]
Pt/ β -SiC HS	open-mouthed hollow spheres	CH ₄	300 W Xe lamp	H ₂ O	-	16.8	[209]
Co ₉ S ₈ @ZnIn ₂ S ₄ /CdS	hollow core-shell nanoreactor	CO	300 W Xe lamp	H ₂ O	-	82.1	[210]
SnS ₂ /CeO ₂	2D/0D heterojunction	CO	300 W Xe lamp	TEOA	-	16.27	[211]
SiW ₁₁ Cu@TiO ₂	heterointerface core-shell-like composite	CO	300 W Xe lamp	TEOA	-	84.12	[212]
Au/CdS HMCHPs	hierarchical multi-cavity hollow particles	CO	300 W Xe lamp $\lambda \geq 400$ nm	TEOA	0.61 (420 nm)	3,758	[213]
Ni/Bi-30	porous ultrathin nanosheets	CO	460 nm LED	TEOA	-	10.2	[214]
NFMS	2D/2D heterojunction	CO	300 W Xe lamp	H ₂ O	-	2.68	[187]

AQE: Apparent quantum efficiency; UN: ultrathin nanosheet; CNGA: glucose-assisted ultrathin porous nitrogen-vacancy carbon nitride nanosheet; CD: carbon dot; CQD: carbon quantum dot; LDH: layered double hydroxide; BMO: Bi₂MoO₆; BWO: Bi₂WO₆; HMCHP: hierarchical multi-cavity hollow particle; NFMS: NiFe-LDH/MoS₂; COF: covalent organic framework; TEOA: triethanolamine; LED: light-emitting diode.

diverse spatiotemporal scales. For example, coupling defect engineering with built-in electric fields, as demonstrated in oxygen-vacancy-rich Ni₂P₂O₇/g-C₃N₄ heterojunctions^[215] or ferroelectric Bi₃TiNbO₉ nanosheets^[216], establishes a powerful dual-regulation mechanism. In these systems, vacancies serve as localized reactant-trapping and activation loci, while the overarching internal polarization fields forcefully dictate directional carrier migration, synchronously suppressing bulk recombination and accelerating surface reaction kinetics.

Furthermore, integrating Schottky junctions with plasmonic architectures introduces distinct energetic advantages. In plasmonic Ag/TiO₂ systems, the localized surface plasmon resonance (LSPR) effect harvests lower-energy photons to inject hot electrons, while the proximal Schottky barrier immediately rectifies their flow, preventing back-transfer. Time-resolved spectroscopy confirms that this precise energetic and spatial synchronization profoundly amplifies overall multi-electron transfer efficiency^[217].

Ultimately, the zenith of photocatalytic CO₂ reduction will not be reached through isolated structural tweaks, but through the mastery of multi-level coupled mechanisms. By hierarchically integrating interfacial electric fields, precise defect distributions, and optimally configured active sites, researchers can construct continuous, energetically cascaded electron “highways” from bulk generation centers directly to adsorbed CO₂ molecules. This synergy-driven design paradigm stands as the definitive roadmap for circumventing intrinsic thermodynamic limitations and steering PCR technology toward viable practical applications.

CONCLUSIONS AND OUTLOOK

This review systematically discusses the charge separation process in photocatalytic CO₂ reduction (PCR) systems from a multi-scale perspective, integrating bulk generation, interfacial migration, and surface reaction coupling within a unified framework. As emphasized throughout the manuscript, charge separation efficiency is the core bottleneck governing PCR performance, as it directly determines the effective utilization of photogenerated carriers in multi-electron CO₂ conversion processes.

By correlating thermodynamic driving forces with kinetic competition, we clarify that charge separation is not an isolated step but a dynamic process governed by the interplay between carrier generation, migration, and recombination across different spatial scales. Intrinsic limitations such as bulk recombination, interfacial energy mismatch, and surface charge consumption inefficiency, together with non-intrinsic factors including light harvesting and CO₂ activation barriers, collectively define the upper limit of photocatalytic efficiency.

From a materials design perspective, we further summarize that charge separation regulation must rely on synergistic strategies rather than single-factor optimization. In particular, heterojunction engineering, including type-II, Z-scheme, S-scheme, and Schottky junctions, provides effective pathways for directional charge transfer and redox potential preservation. Among these, Z-scheme and S-scheme systems are especially effective in balancing charge separation with strong redox capability, while Schottky junctions offer an additional route for surface electron extraction and reaction-site coupling.

In addition, as highlighted in the characterization section, reliable evaluation of charge separation requires a multi-technique and multi-timescale approach. No single method can fully describe the complete carrier evolution process; instead, only the convergence of spectroscopic, electrochemical, ultrafast kinetic, and operando techniques can reconstruct the full pathway from photogeneration to catalytic consumption. Overall, this work establishes a unified picture linking charge generation, separation, transport, and utilization in PCR systems, providing mechanistic guidance for the rational design of high-efficiency photocatalysts.

Despite the remarkable progress in regulating charge separation for photocatalytic CO₂ reduction (PCR), several critical challenges remain before translating these advances into practical applications. Future research should move beyond isolated optimization strategies toward data-driven, mechanism-oriented, and application-relevant catalyst design frameworks, with particular emphasis on charge behavior under realistic reaction environments.

First, AI-assisted catalyst design is expected to play a transformative role in accelerating the discovery of high-performance PCR systems. Machine learning has been widely recognized as an effective tool for accelerating molecular and materials discovery, and its application in electrocatalyst and photocatalyst design has provided new opportunities for identifying structure-property-activity relationships^[218,219]. By integrating machine learning with density functional theory calculations and high-throughput experimentation, it becomes possible to establish quantitative correlations between electronic structure descriptors, such as band alignment, charge density distribution, carrier lifetime, and catalytic performance. In particular, AI models can be trained to predict charge separation efficiency, interfacial charge-transfer kinetics, and optimal heterojunction configurations, thereby guiding the rational design of multi-component catalysts beyond empirical trial-and-error approaches.

Second, the development of advanced operando characterization techniques is essential to directly probe charge behavior under working conditions. Although PL, TRPL, EIS, and TAS provide valuable information on carrier recombination and transport, they often cannot fully capture real-time charge redistribution during CO₂ adsorption, activation, and conversion. Emerging *in situ*/operando spectroscopies, including operando XPS, XAS, FT-IR/DRIFTS, Raman, and EPR, can provide more direct evidence for photoinduced charge migration, interfacial band bending, active-site evolution, and reaction intermediate formation under illumination^[220]. Therefore, future studies should combine ultrafast spectroscopy, operando electronic-structure characterization, and theoretical calculations to reconstruct the full charge evolution pathway from photogeneration to surface reaction.

Third, the rational design of multi-strategy synergistic catalysts should be significantly strengthened. Instead of relying on a single modulation strategy, future PCR catalysts should integrate heterojunction engineering, defect regulation, cocatalyst loading, morphology optimization, and internal electric-field construction into unified architectures. In such systems, charge separation is not governed by one isolated interface but by multi-level coupling across the bulk, interface, and surface. For example, ferroelectric polarization coupled with oxygen vacancies can simultaneously promote bulk charge separation and surface CO₂ activation^[220], while plasmonic metal/semiconductor interfaces can couple Schottky-barrier-driven charge extraction with hot-electron injection and enhanced light absorption. These examples indicate that synergistic catalyst design should aim to regulate the entire charge-transfer network, including carrier generation, directional migration, interfacial accumulation, and final utilization at CO₂ reduction sites.

Fourth, a deeper understanding of charge separation mechanisms under conditions closer to practical application is urgently needed. Most current PCR studies are still performed under idealized conditions, such as high-purity CO₂, sacrificial agents, static batch reactors, and simplified gas-solid or liquid-solid interfaces. However, under realistic conditions, including low-concentration CO₂, fluctuating light intensity, water vapor, gas-liquid-solid interfaces, and continuous-flow operation, charge behavior may be strongly affected by mass transfer, competitive adsorption, surface reconstruction, and local reaction microenvironments^[221,222]. Therefore, future research should investigate charge transport, recombination, and interfacial electron utilization in diluted CO₂ systems and continuous-flow photoreactors. Such studies will be essential for bridging the gap between laboratory-scale photocatalytic activity and scalable CO₂ conversion systems.

Finally, from a fundamental perspective, future research should aim to establish a unified theoretical framework for charge behavior that integrates thermodynamics, kinetics, and interfacial physics. This framework should quantify the competition between charge transfer and recombination across multiple spatial and temporal scales and identify key descriptors governing charge separation efficiency. In summary, advancing photocatalytic CO₂ reduction toward practical implementation requires a paradigm shift from

material-centered optimization to charge-behavior-centered design, supported by AI-driven catalyst discovery, operando characterization, multi-strategy synergistic engineering, and realistic-condition mechanism studies. These directions will provide new opportunities for developing highly efficient, selective, stable, and scalable photocatalytic CO₂ reduction systems.

DECLARATIONS

Authors' contributions

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During the preparation of this manuscript, the AI tool ChatGPT (version 5.5, released 2026-04-24) was used solely for language editing. The tool did not influence the study design, data collection, analysis, interpretation, or the scientific content of the work. All authors take full responsibility for the accuracy, integrity, and final content of the manuscript.

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Conflicts of interest

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Ethical approval and consent to participate

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